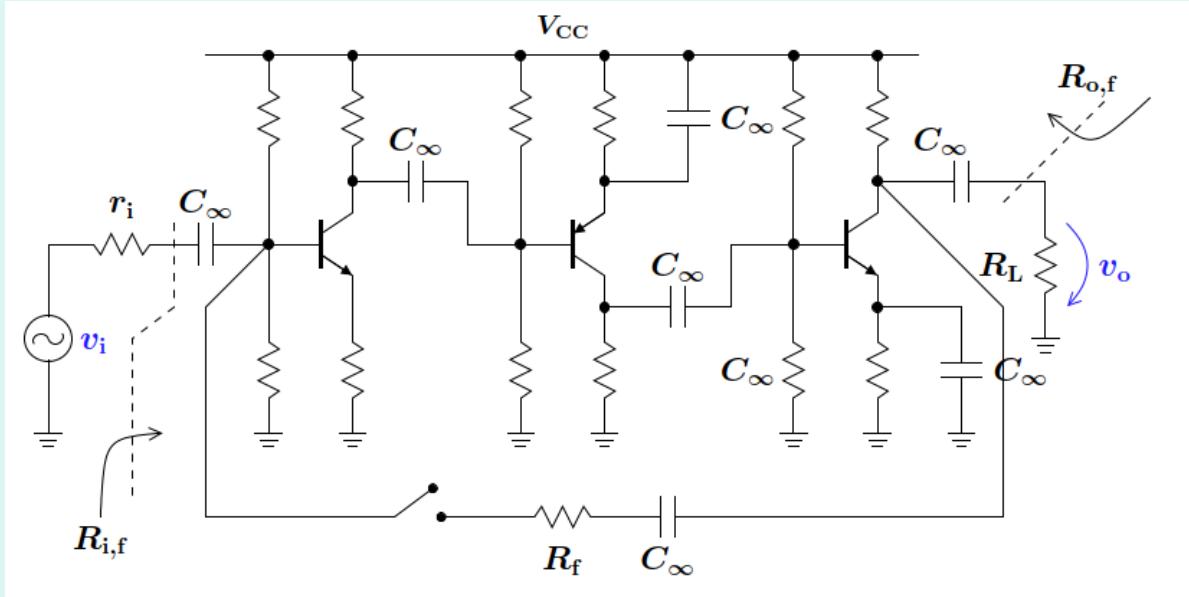
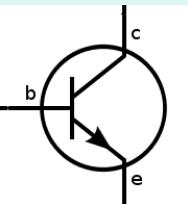


# ELEC 301 - Band-pass circuits

L08 - Sep 22

Instructor: Edmond Cretu

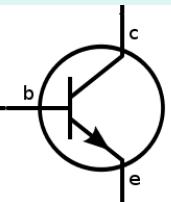




# Last time

- Bode plots - log magnitude and phase plots
- Examples
- Approximation errors

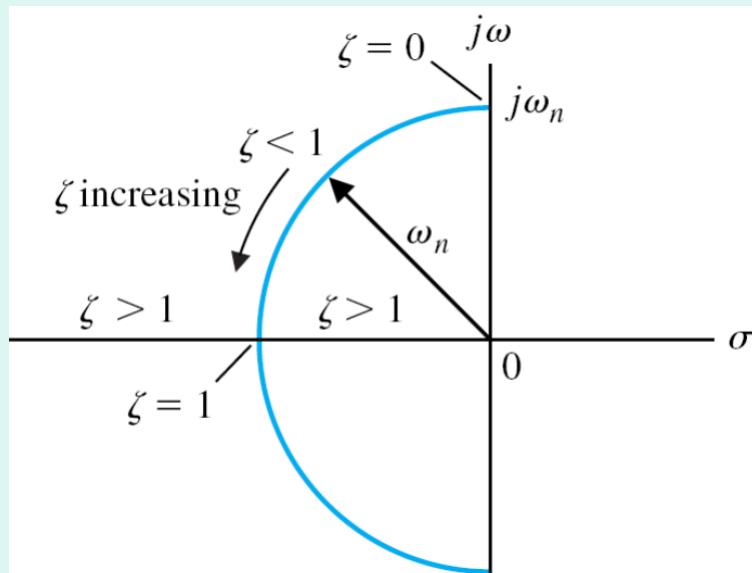


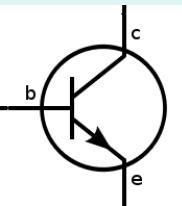


# L08 Q01 - Bode plots for a 2<sup>nd</sup> order system

- Given a second order system, in which case the Bode technique provides a better approximation of the real frequency response?
  - quality factor  $Q=1000$
  - quality factor  $Q=10$
  - Quality factor  $Q=0.5$

$$Y(s) = \frac{\omega_n^2}{s^2 + 2\zeta\omega_n s + \omega_n^2} X(s)$$

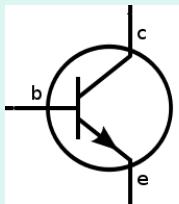




# Conclusions

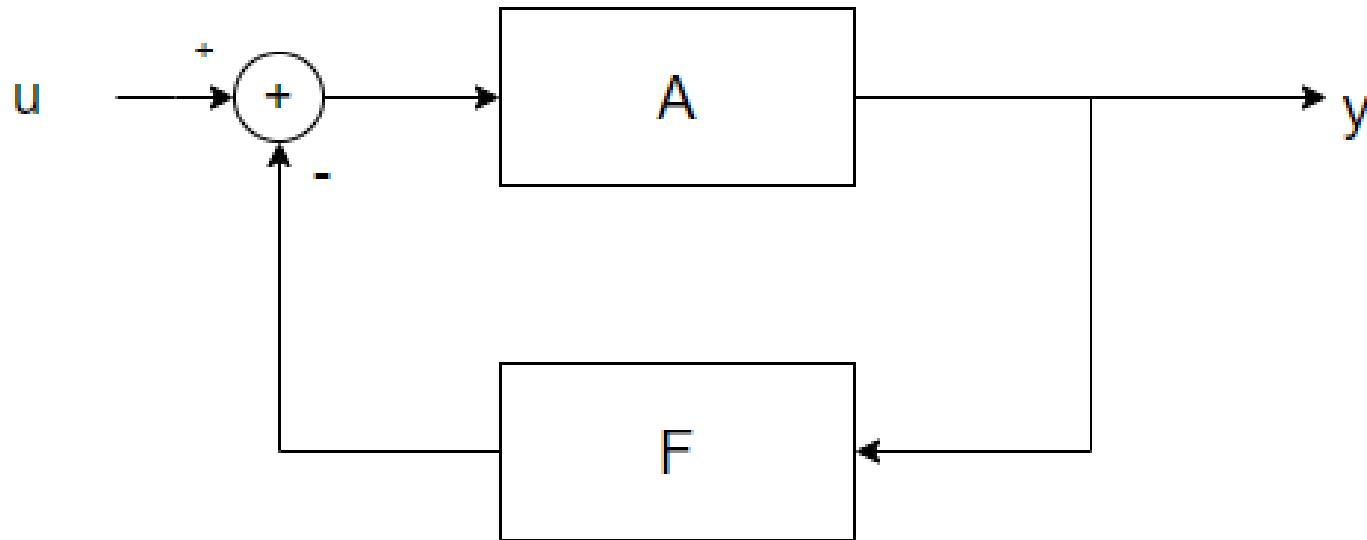
- Bode plots is a fast visual approximation technique for the frequency response of a linear system
- It provides a good approximation for systems with distinct individual poles on the real axis
- The typical cases for electronic amplifiers correspond to such a situation





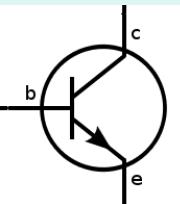
# Simple LP active amplifier

- Simple low-pass amplifier stage with passive feedback network



$$A(j\omega) = \frac{A_0}{1 + j \frac{\omega}{\omega_H}}, \quad F = F_0 = \frac{R_1}{R_1 + R_2} < 1$$

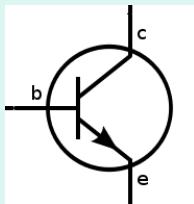




## L08 Q02 LP amplifier

- Which statement is **false** for the previous configuration?
  - A. The closed loop gain decreases because of the negative feedback
  - B. The critical frequency of the closed loop transfer function decreases as result of the negative feedback action
  - C. The critical frequency of the closed loop transfer function increases, compared to the open loop case
  - D. The closed loop system is less sensitive to errors in the value of  $A_0$



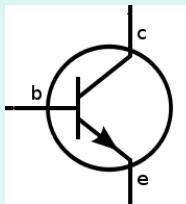


# Negative feedback action

- The critical frequency  $\omega_H$  is increased by the loop gain

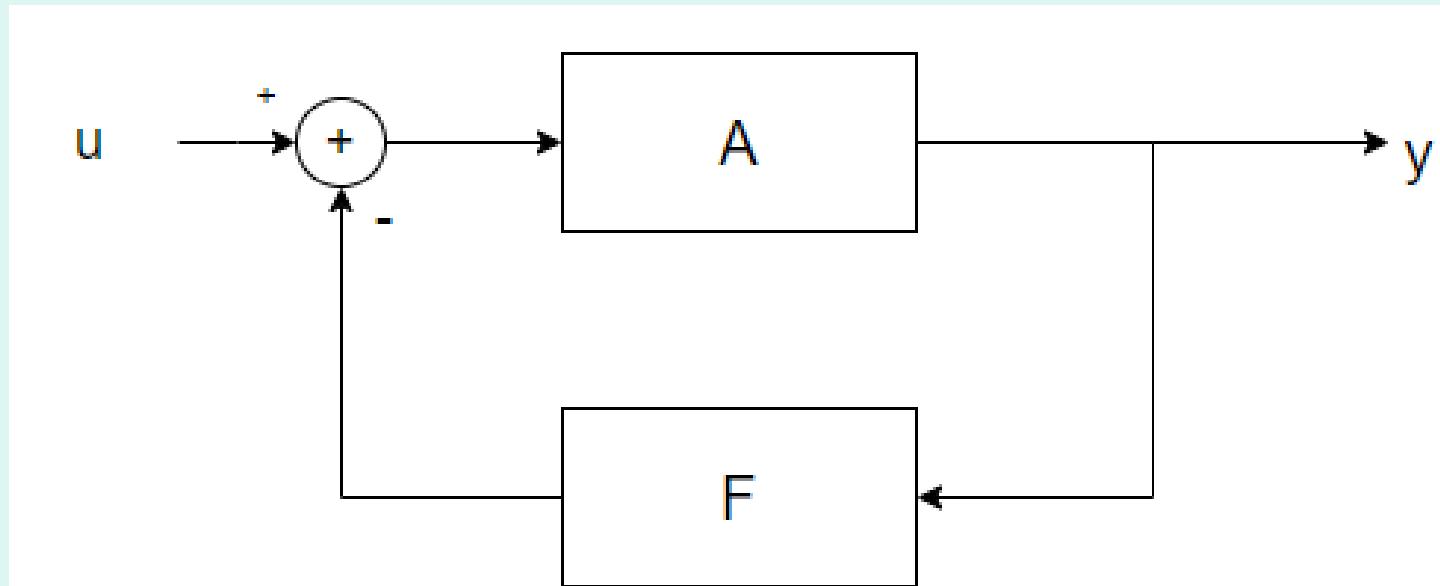
$$\begin{aligned}
 A_f(j\omega) &= \frac{A(j\omega)}{1 + A(j\omega)F(j\omega)} = \frac{\frac{A_0}{1 + j\frac{\omega}{\omega_H}}}{1 + \frac{A_0}{1 + j\frac{\omega}{\omega_H}}F_0} = \frac{A_0}{1 + j\frac{\omega}{\omega_H} + A_0F_0} \\
 A_f(j\omega) &= \frac{A_0}{1 + A_0F_0} \frac{1}{1 + j\frac{\omega}{\omega_H(1 + A_0F_0)}} \underset{A_0F_0 \gg 1}{\approx} \frac{1}{F_0} \frac{1}{1 + j\frac{\omega}{\omega_H(1 + A_0F_0)}}
 \end{aligned}$$





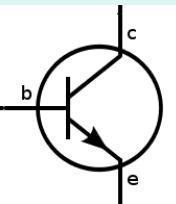
# High-pass amplifier case

- Similar active network, but with a high-pass amplifier characteristics



$$A(j\omega) = \frac{A_0 j \frac{\omega}{\omega_L}}{1 + j \frac{\omega}{\omega_L}}, \quad F = F_0 = \frac{R_1}{R_1 + R_2} < 1$$

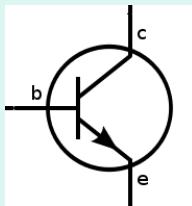




# L08 Q03 HP amplifier with feedback

- What is the effect of the negative feedback on the critical frequency  $\omega_L$ ?
  - A. It is not changed
  - B. It is increased
  - C. It is decreased





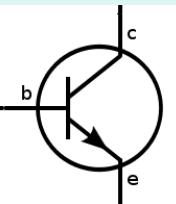
# HP active amplifier with feedback

- The critical frequency is lowered by the loop gain

$$A(j\omega) = \frac{A_0 j \frac{\omega}{\omega_L}}{1 + j \frac{\omega}{\omega_L}}, \quad F = F_0 = \frac{R_1}{R_1 + R_2} < 1$$

$$_c(j\omega) = \frac{A_0 j \frac{\omega}{\omega_L}}{A_0 F_0 j \frac{\omega}{\omega_L} + 1 + j \frac{\omega}{\omega_L}} = \frac{A_0}{1 + A_0 F_0} \frac{j \frac{\omega}{\omega_L}}{1 + j \frac{\omega}{\omega_L} \frac{(1 + A_0 F_0)}{(1 + A_0 F_0)}}$$

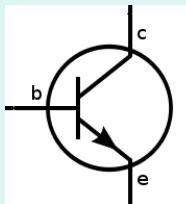




# Open-circuit/short-circuit time constant method

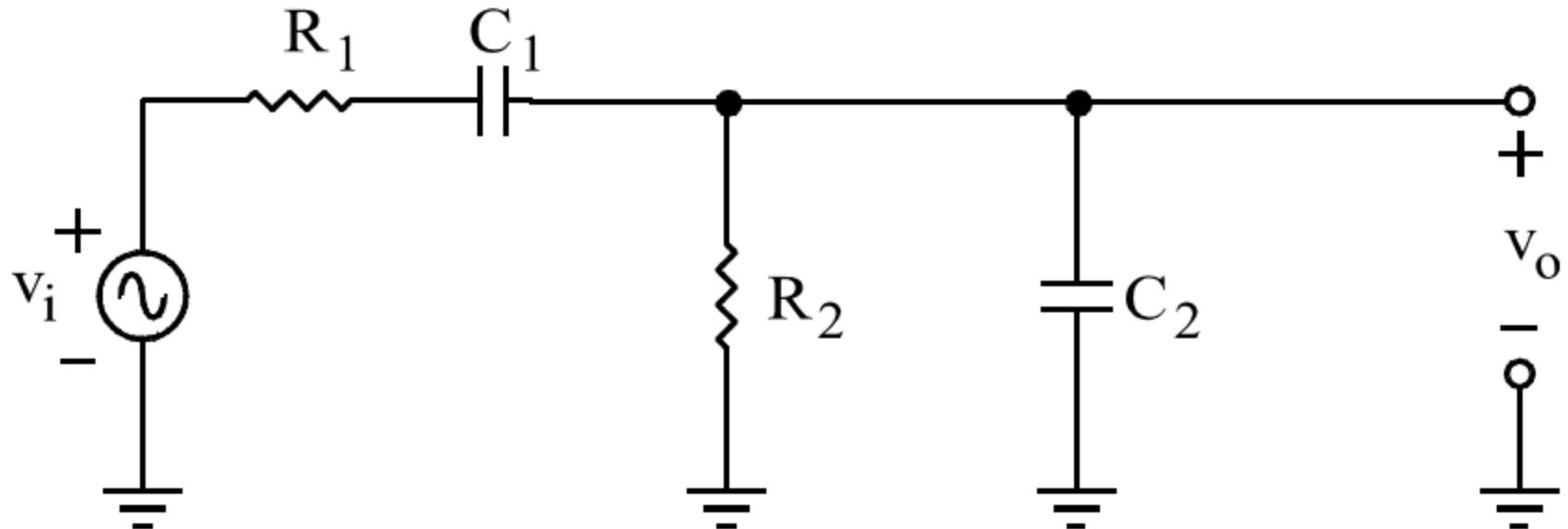
- Most circuits we operate with are bandpass, with capacitors as the only reactive components
- We desire a fast method to approximate the critical frequencies  $\omega_L$ ,  $\omega_H$

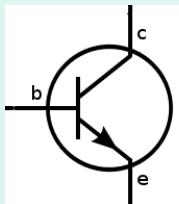




# Example - simple bandpass filter

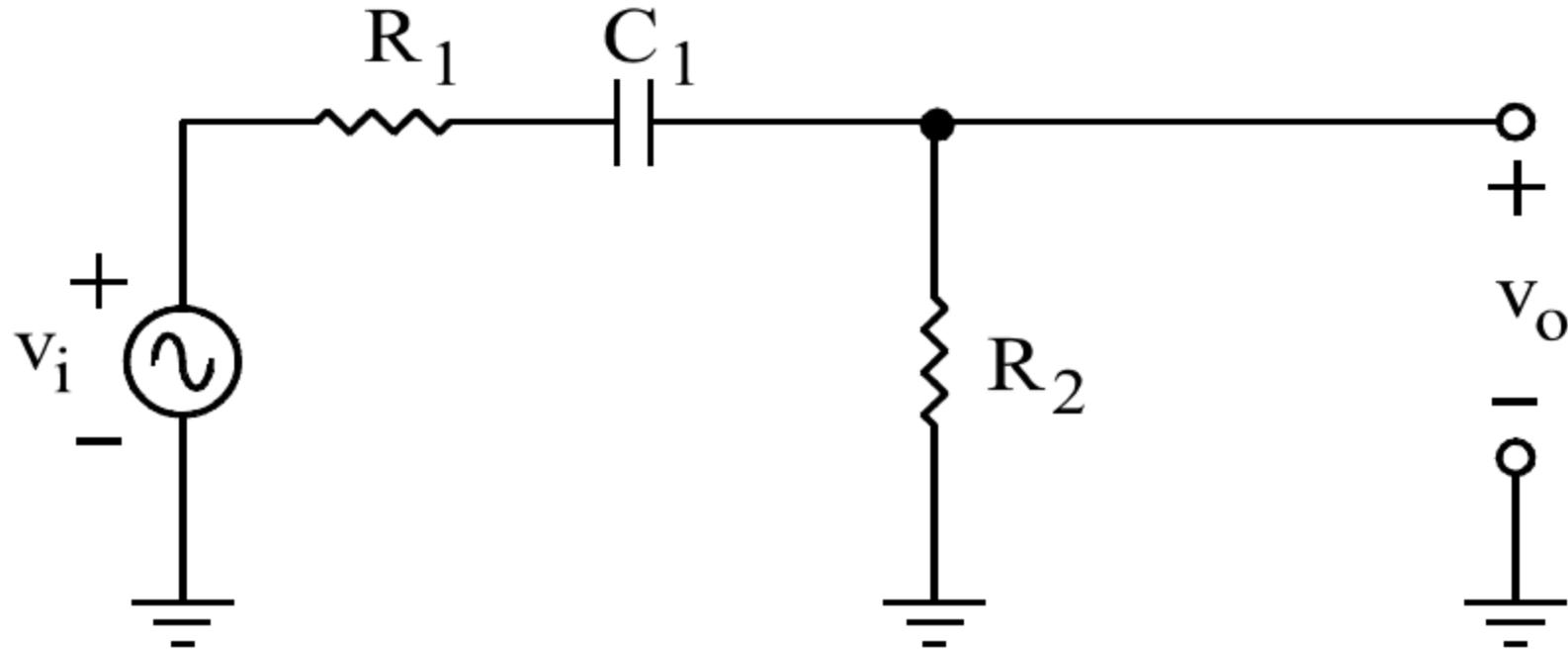
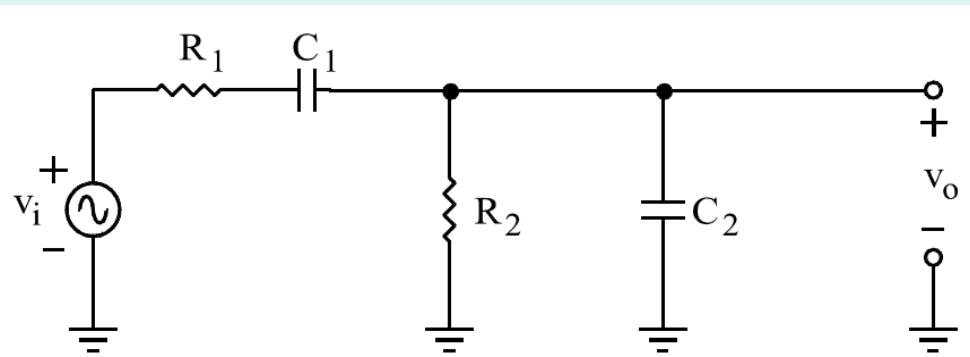
- $C_1$  blocks the LF components,  $C_2$  attenuates the HF components
- To be a bandpass filter  $C_1 \gg C_2$

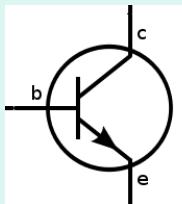




# Circuit at LF

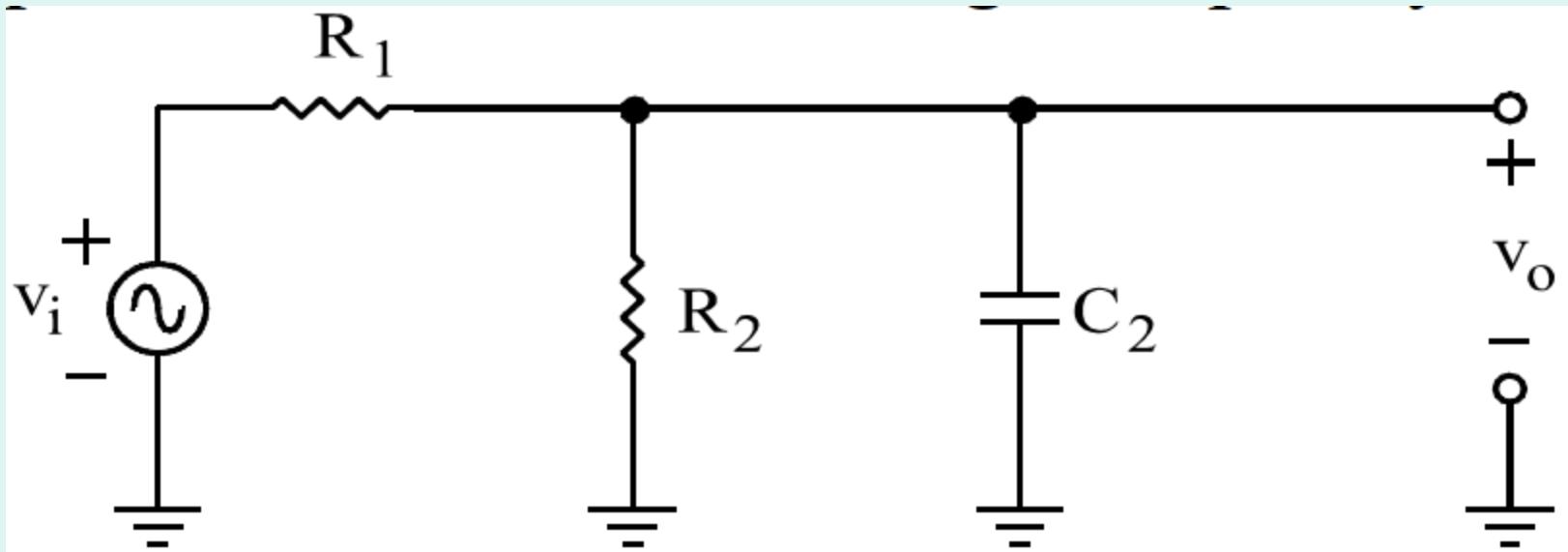
- Assume  $C_1 \gg C_2$  - circuit behaves as high-pass filter for low frequency range

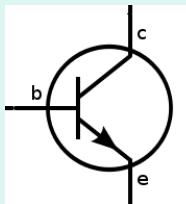




# Circuit at HF

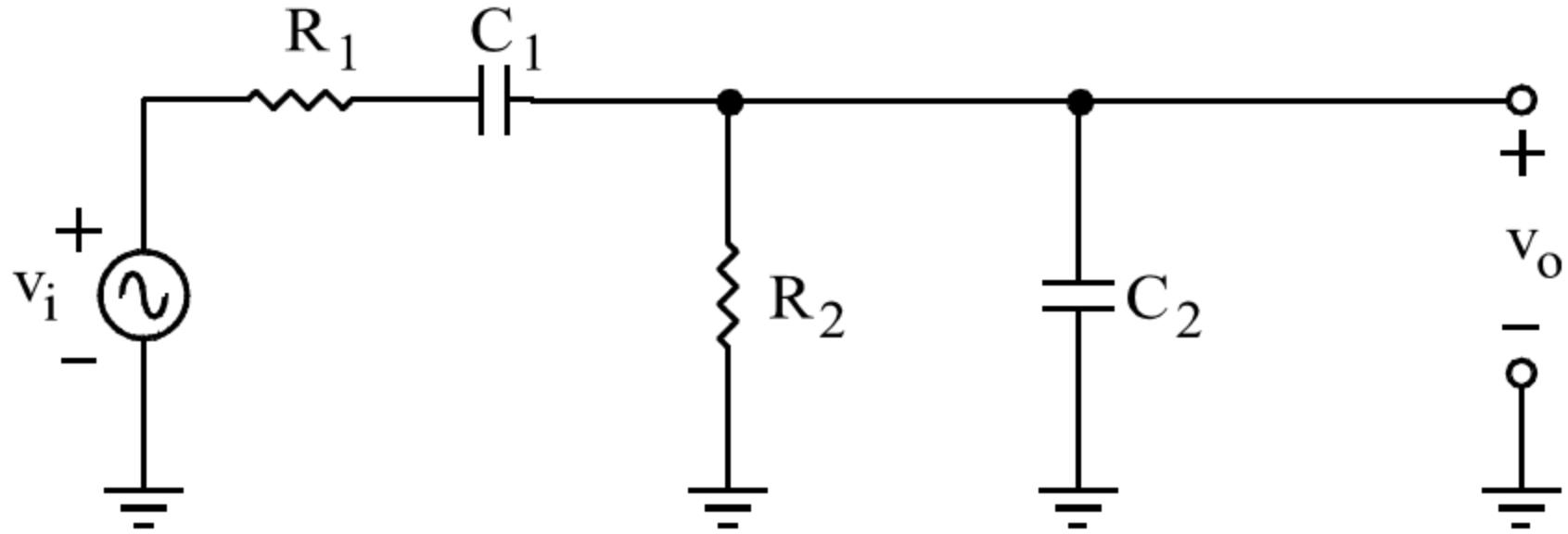
- Assume  $C_1 \gg C_2$  - circuit behaves as low-pass filter for high frequency range





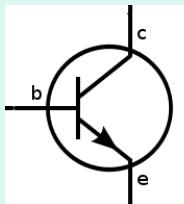
# Transfer function

- Homework!



$$\frac{V_o(s)}{V_i(s)} = \frac{\frac{s}{R_1 C_2}}{s^2 + \left[ \frac{1}{R_2 C_2} + \frac{1}{R_1 C_2} + \frac{1}{R_1 C_1} \right] s + \frac{1}{R_1 R_2 C_1 C_2}}$$

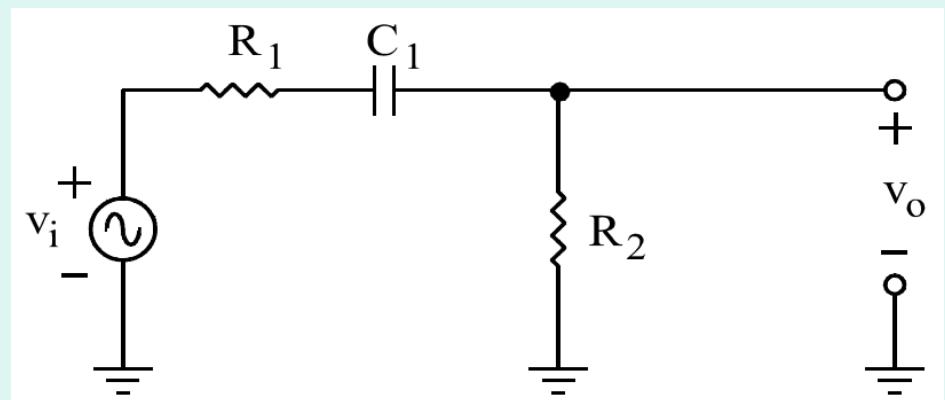




# Approximate LF pole location

- Recall the circuit approximations for the LF and HF cases

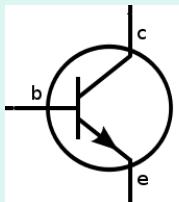
LF regime



$$\left( \frac{V_o}{V_i} \right)_{LF} = \frac{R_2}{R_2 + R_1 + \frac{1}{sC_1}} = \frac{sR_2C_1}{1 + s(R_1 + R_2)C_1}$$

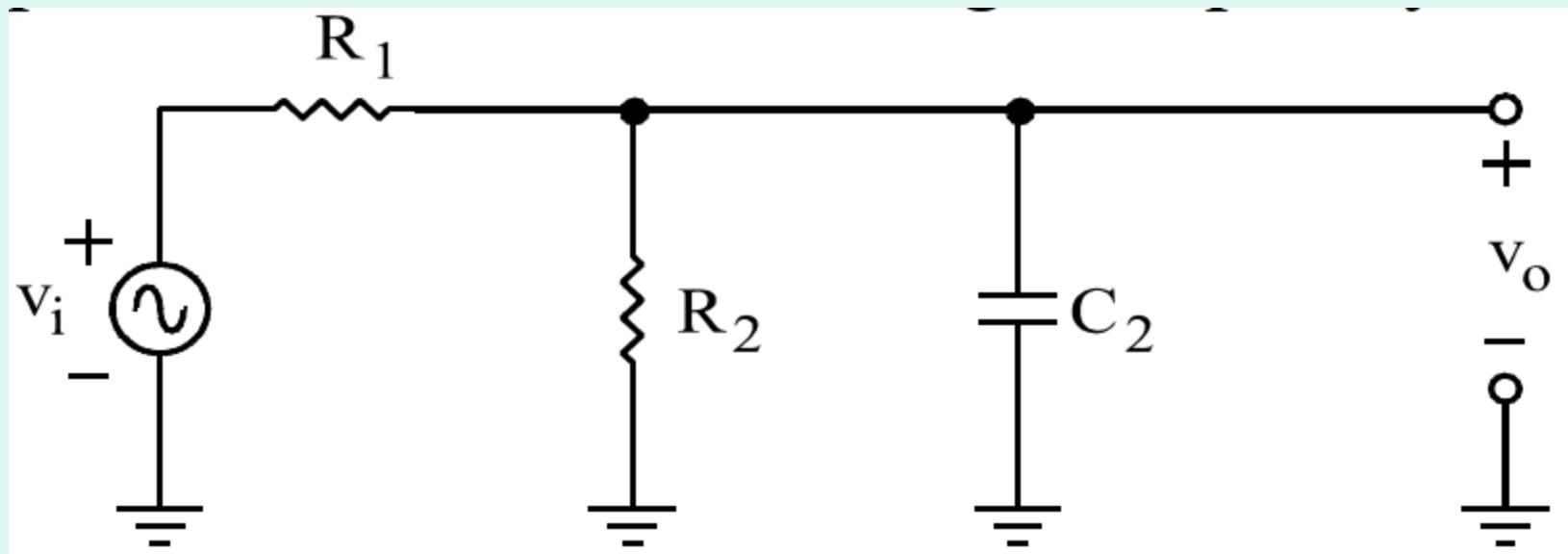
$$\omega_{p1} \approx \frac{1}{(R_1 + R_2)C_1} = \frac{1}{R_+ C_1}$$





# Approximate HF pole location

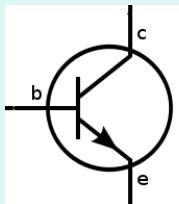
- The simplified circuit at HF:



$$\left(\frac{V_o}{V_i}\right)_{HF} \approx \frac{R_2 // \frac{1}{sC_2}}{R_1 + R_2 // \frac{1}{sC_2}} = \frac{\frac{R_2}{1 + sR_2C_2}}{R_1 + \frac{R_2}{1 + sR_2C_2}} = \frac{R_2}{R_1 + R_2 + sR_1R_2C_2}$$

  $\left(\frac{V_o}{V_i}\right)_{HF} \approx \frac{R_2}{R_1 + R_2} \frac{1}{1 + sC_2 \left( \frac{1}{R_1} + \frac{1}{R_2} \right)^{-1}}$

$$\omega_{p2} \approx \frac{1}{\left( \frac{1}{R_1} + \frac{1}{R_2} \right)^{-1} C_2} = \frac{1}{(R_1 // R_2) C_2}$$



# Combined frequency response

- With the obtained poles approximations:

$$\frac{V_o(s)}{V_i(s)} = \frac{\frac{s}{R_1 C_2}}{s^2 + \left[ \frac{1}{R_2 C_2} + \frac{1}{R_1 C_2} + \frac{1}{R_1 C_1} \right] s + \frac{1}{R_1 R_2 C_1 C_2}}$$

$$\frac{V_o(s)}{V_i(s)} = \frac{\frac{s}{R_1 C_2}}{\left( s + \frac{1}{R_+ C_1} \right) \left( s + \frac{1}{R_{||} C_2} \right)} = \frac{R_2}{R_1 + R_2} \frac{s}{\left( s + \frac{1}{R_+ C_1} \right) \left( s + \frac{1}{R_{||} C_2} \right)}$$

$$= \frac{R_2}{R_1 + R_2} \frac{s}{(s + \omega_{p1}) (s + \omega_{p2})}$$

↗

↑

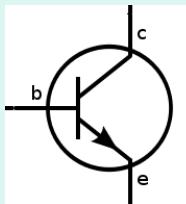
↖

Midband

Low freq.

High freq.





# Generalization

- We can in general, provided that the poles groups are apart, separate an amplifier frequency response into the three bands components:

$$T(s) = A_M F_L(s) F_H(s)$$

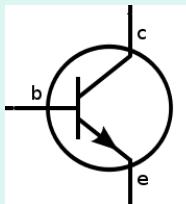
↗      ↑      ↘

Midband      Lowf      Highf  
(high-pass)   (low-pass)

$$F_L(s) = \frac{(s + \omega_{Lz1}) \cdots (s + \omega_{Lzn})}{(s + \omega_{Lp1}) \cdots (s + \omega_{LpN})}$$



Here  $n=N$  to ensure that  $|F_L(\infty)|=1$



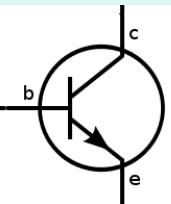
## Generalization (2)

- For the HF component:

$$\begin{aligned} F_H(s) &= \frac{\omega_{Hp1} \cdots \omega_{HpM}}{\omega_{Hz1} \cdots \omega_{Hzm}} \frac{(s + \omega_{Hz1}) \cdots (s + \omega_{Hzm})}{(s + \omega_{Hp1}) \cdots (s + \omega_{HpM})} \\ &= \frac{\left(1 + \frac{s}{\omega_{Hz1}}\right) \cdots \left(1 + \frac{s}{\omega_{Hzm}}\right)}{\left(1 + \frac{s}{\omega_{Hp1}}\right) \cdots \left(1 + \frac{s}{\omega_{HpM}}\right)} \end{aligned}$$

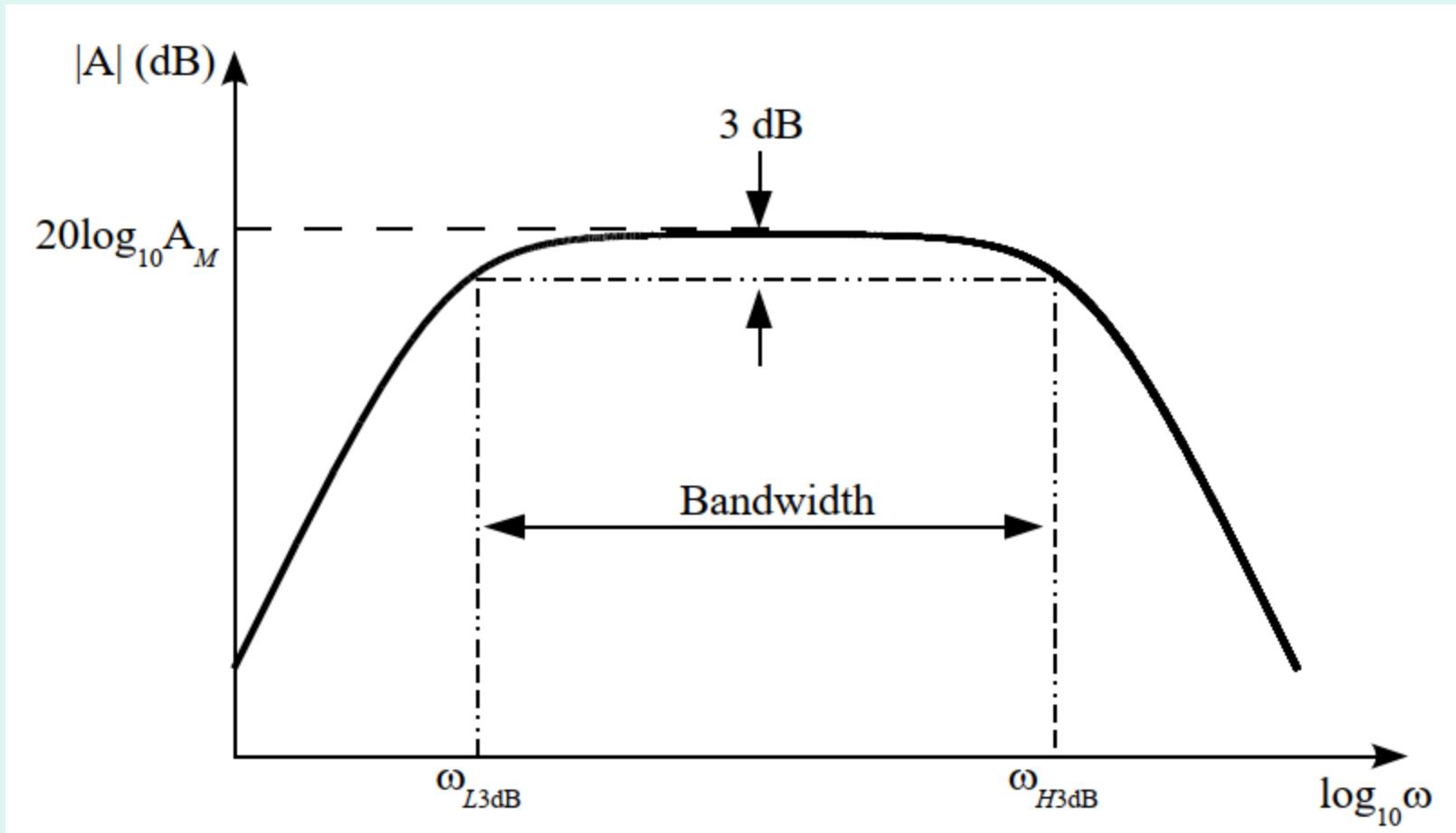
Here  $M > m$ , to ensure that  $|F_H(\infty)| = 0$

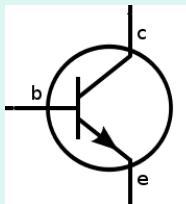




# Generic amplifier frequency response

- Bandwidth  $BW = \omega_{H3dB} - \omega_{L3dB}$





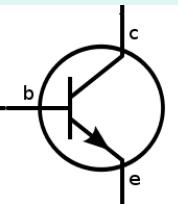
# Finding the cut-off frequency $\omega_{L3dB}$

- We assume we know the transfer functions (locations of poles and zeros)

$$|F_L(j\omega_{L3dB})|^2 = \frac{(\omega_{L3dB}^2 + \omega_{Lz1}^2)(\omega_{L3dB}^2 + \omega_{Lz2}^2) \cdots (\omega_{L3dB}^2 + \omega_{Lzn}^2)}{(\omega_{L3dB}^2 + \omega_{Lp1}^2)(\omega_{L3dB}^2 + \omega_{Lp2}^2) \cdots (\omega_{L3dB}^2 + \omega_{LpN}^2)} = \frac{1}{2}$$

$$2(\omega_{L3dB}^2 + \omega_{Lz1}^2) \cdots (\omega_{L3dB}^2 + \omega_{Lzn}^2) = (\omega_{L3dB}^2 + \omega_{Lp1}^2) \cdots (\omega_{L3dB}^2 + \omega_{LpN}^2)$$

$$\begin{aligned} & 2(\omega_{L3dB}^{2n} + (\omega_{Lz1}^2 + \omega_{Lz2}^2 + \cdots + \omega_{Lzn}^2)\omega_{L3dB}^{2(n-1)} + \cdots + (\omega_{Lz1}^2 \omega_{Lz2}^2 \cdots \omega_{Lzn}^2)) \\ & = (\omega_{L3dB}^{2N} + (\omega_{Lp1}^2 + \omega_{Lp2}^2 + \cdots + \omega_{LpN}^2)\omega_{L3dB}^{2(N-1)} + \cdots + (\omega_{Lp1}^2 \omega_{Lp2}^2 \cdots \omega_{LpN}^2)) \end{aligned}$$

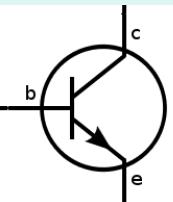


## Finding the cut-off frequency $\omega_{L3dB}$ (2)

- Since  $\omega_{L3dB}$  is larger than any pole or zero, we can approximate the eqn. considering only the highest powers
- Rem: many of the zeros may be located at 0

$$2\omega_{L3dB}^{2N} + 2(\omega_{Lz1}^2 + \dots + \omega_{Lzn}^2)\omega_{L3dB}^{2(N-1)} \approx \omega_{L3dB}^{2N} + (\omega_{Lp1}^2 + \dots + \omega_{LpN}^2)\omega_{L3dB}^{2(N-1)}$$

$$\omega_{L3dB} \approx \sqrt{\omega_{Lp1}^2 + \dots + \omega_{LpN}^2 - 2\omega_{Lz1}^2 - \dots - 2\omega_{Lzn}^2}$$

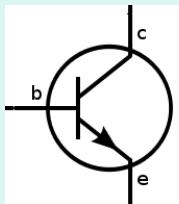


# Finding the cut-off frequency $\omega_{3\text{dBH}}$

- At  $\omega_{3\text{dBH}}$ , the normalized  $F_H(s)$  should be reduced by  $\sqrt{2}$

$$|F_H(j\omega_{3\text{dB}})|^2 = \frac{\left(1 + \frac{\omega_{H3\text{dB}}^2}{\omega_{Hz1}^2}\right) \left(1 + \frac{\omega_{H3\text{dB}}^2}{\omega_{Hz2}^2}\right) \cdots \left(1 + \frac{\omega_{H3\text{dB}}^2}{\omega_{Hzm}^2}\right)}{\left(1 + \frac{\omega_{H3\text{dB}}^2}{\omega_{Hp1}^2}\right) \left(1 + \frac{\omega_{H3\text{dB}}^2}{\omega_{Hp2}^2}\right) \cdots \left(1 + \frac{\omega_{H3\text{dB}}^2}{\omega_{HpM}^2}\right)} = \frac{1}{2}$$

$$2 \left(1 + \frac{\omega_{H3\text{dB}}^2}{\omega_{Hz1}^2}\right) \left(1 + \frac{\omega_{H3\text{dB}}^2}{\omega_{Hz2}^2}\right) \cdots \left(1 + \frac{\omega_{H3\text{dB}}^2}{\omega_{Hzm}^2}\right) = \left(1 + \frac{\omega_{H3\text{dB}}^2}{\omega_{Hp1}^2}\right) \left(1 + \frac{\omega_{H3\text{dB}}^2}{\omega_{Hp2}^2}\right) \cdots \left(1 + \frac{\omega_{H3\text{dB}}^2}{\omega_{HpM}^2}\right)$$



# Finding the cut-off frequency $\omega_{3\text{dBH}}^2$ (2)

- Since  $M > m$  and  $\omega_{3\text{dBH}}^2$  is smaller than any of the pole or zero frequencies, we can approximate the equality considering only the (dominant) low powers in the polynomials

$$2 + \left( \frac{2}{\omega_{Hz1}^2} + \dots + \frac{2}{\omega_{Hzm}^2} \right) \omega_{H3\text{dB}}^2 \approx 1 + \left( \frac{1}{\omega_{Hp1}^2} + \dots + \frac{1}{\omega_{HpM}^2} \right) \omega_{H3\text{dB}}^2$$

$$\frac{1}{\omega_{H3\text{dB}}^2} \approx \frac{1}{\omega_{Hp1}^2} + \dots + \frac{1}{\omega_{HpM}^2} - \frac{2}{\omega_{Hz1}^2} - \dots - \frac{2}{\omega_{Hzm}^2}$$

$$\tau_{H3\text{dB}}^2 \approx \sqrt{\tau_{Hp1}^2 + \dots + \tau_{HpM}^2 - 2\tau_{Hz1}^2 - \dots - 2\tau_{Hzm}^2}$$

$$\omega_{H3\text{dB}} = \frac{1}{\tau_{H3\text{dB}}}$$

