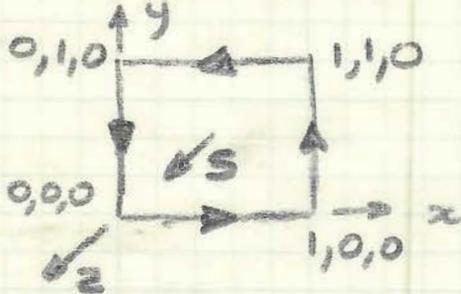


Given $\mathbf{H} = 300\mathbf{a}_z \cos(3 \times 10^8 t - y)$ A/m in free space, find the emf developed in the general \mathbf{a}_ϕ direction about the closed path having corners at (a) (0, 0, 0), (1, 0, 0), (1, 1, 0), and (0, 1, 0); (b) (0, 0, 0) (2 π , 0, 0), (2 π , 2 π , 0), and (0, 2 π , 0).

a.

• THE PATH LIES WITHIN THE x-y PLANE



• THE MAGNETIC FLUX DENSITY IS ALIGNED WITH THE NORMAL TO THE PATH

$$\mathbf{B} = \mu_0 \mathbf{H} = 300\mu_0 \cos(3 \times 10^8 t - y) \hat{\mathbf{z}} \text{ A/m}$$

$$\text{emf} = \frac{d\Phi}{dt} \quad \Phi = \iint_S \bar{\mathbf{B}} \cdot d\bar{\mathbf{S}}$$

$$\Phi = \int_0^1 \int_0^1 300\mu_0 \cos(3 \times 10^8 t - y) dx dy$$

$$= 300\mu_0 \sin(3 \times 10^8 t - y) \Big|_0^1$$

$$= 300\mu_0 [\sin(3 \times 10^8 t - 1) - \sin(3 \times 10^8 t)] \text{ wb}$$

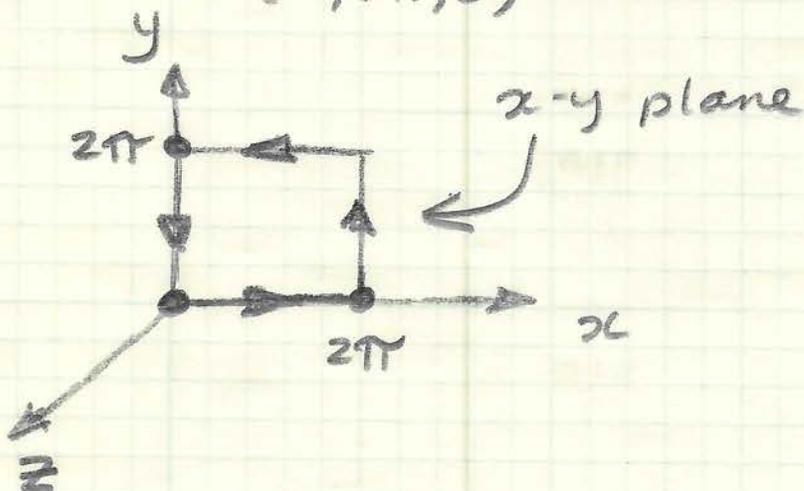
$$\text{emf} = -\frac{d\Phi}{dt}$$

$$= 300(3 \times 10^8)(4\pi \times 10^{-7}) [\cos(3 \times 10^8 t - 1) - \cos(3 \times 10^8 t)]$$

$$= -1.13 \times 10^5 [\cos(3 \times 10^8 t - 1) - \cos(3 \times 10^8 t)] \text{ V}$$

b. CORNERS AT

$(0,0,0)$ $(2\pi,0,0)$ $(2\pi,2\pi,0)$
 $(0,2\pi,0)$



$$\Phi(t) = 2\pi \cdot 300 \mu_0 \sin(3 \times 10^8 t - y) \Big|_0^{2\pi}$$
$$= 0 \text{ Wb.}$$

emf IS THEREFORE 0 AS WELL.

9.5

9.5 The location of the sliding bar in Figure 9.5 is given by $x = 5t + 2t^3$, and the separation of the two rails is 20 cm. Let $\mathbf{B} = 0.8x^2\mathbf{a}_z$ T. Find the voltmeter reading at (a) $t = 0.4$ s; (b) $x = 0.6$ m.

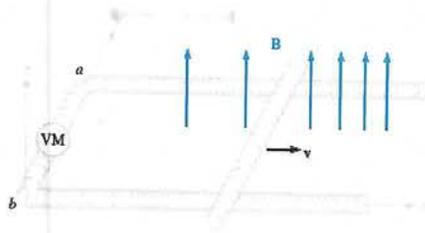


Fig. 9.5

a) $t = 0.4$ s

$$\textcircled{1} \quad \Phi = \int_0^{0.2} \int_0^x 0.8 x'^2 dx' dy \Leftarrow \iint \mathbf{B} \cdot d\mathbf{S}$$

$$= \frac{0.16}{3} x^3 = \frac{0.16}{3} (5t + 2t^3)^3 \text{ Wb}$$

$$\textcircled{2} \quad \text{emf} = -\frac{d\Phi}{dt} = -\frac{0.16}{3} 3(5t + 2t^3)^2 (5 + 6t^2)$$

$$= -0.16 [5 \times 4 + 2 \times 4^3]^2 [5 + 6 \cdot 4^2]$$

$$= -4.32 \text{ V}$$

b) $x = 0.6$ m $\quad 0.6 = 5t + 2t^3 \Rightarrow$
 $t = 0.1193$

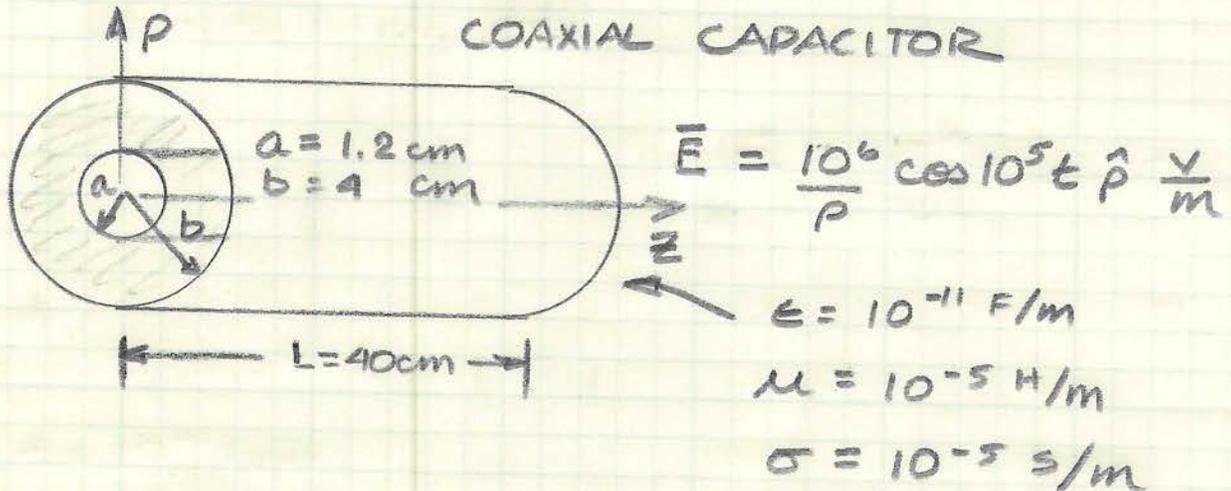
$$\textcircled{2} \quad \text{emf} = -\frac{d\Phi}{dt}$$

$$= -0.16 [5(.1193) + 2(.1193)^3]^2 [5 + 6(.1193)^2]$$

$$= -0.293 \text{ V}$$

9.11

- 9.11 Let the internal dimensions of a coaxial capacitor be $a = 1.2$ cm, $b = 4$ cm, and $l = 40$ cm. The homogeneous material inside the capacitor has the parameters $\epsilon = 10^{-11}$ F/m, $\mu = 10^{-5}$ H/m, and $\sigma = 10^{-5}$ S/m. If the electric field intensity is $\mathbf{E} = (10^6/\rho) \cos 10^5 t \hat{\rho}$ V/m, find (a) \mathbf{J} ; (b) the total conduction current I_c through the capacitor; (c) the total displacement current I_d through the capacitor; (d) the ratio of the amplitude of I_d to that of I_c , the quality factor of the capacitor.



a. $\vec{J} = \sigma \vec{E} = \frac{10}{\rho} \cos 10^5 t \hat{\rho} \text{ A/m}^2$

b. $I_c = \iint \vec{J} \cdot d\vec{S} \quad dS = d\phi dL$

$$= \frac{10}{\rho} \cos 10^5 t \cdot 2\pi \rho L$$

$$= 8\pi \cos 10^5 t \text{ A} = \text{TOTAL CONDUCTION CURRENT}$$

c. $\vec{J}_d = \frac{\partial \vec{D}}{\partial t} = \frac{\partial \epsilon \vec{E}}{\partial t} = \epsilon \frac{\partial \vec{E}}{\partial t}$

$$= \frac{10^{-11} \cdot 10^6 \cdot 10^5}{\rho} \sin 10^5 t \hat{\rho} \frac{A}{m}$$

$$I_d = 2\pi \rho L J_d = -2\pi L \sin 10^5 t$$

$$= -0.8\pi \sin 10^5 t \text{ A} = \text{TOTAL DISPLACEMENT CURRENT}$$

d. $|I_d|/|I_c| = 0.8/8 = 0.1 = \text{QUALITY FACTOR OF THE CAPACITOR}$

9.12 The magnetic flux density $\mathbf{B} = B_0 \cos(\omega t) \cos(k_0 z) \mathbf{a}_y$ Wb/m² exists in free space. B_0 and k_0 are constants. Find (a) the displacement current density; (b) the electric field intensity; (c) k_0 .

a. DISPLACEMENT CURRENT DENSITY

$$\nabla \times \bar{\mathbf{H}} = \frac{\partial \bar{\mathbf{D}}}{\partial t} = \epsilon_0 \frac{\partial \bar{\mathbf{E}}}{\partial t} = \nabla \times \frac{\bar{\mathbf{B}}}{\mu_0} = \bar{\mathbf{J}}_d$$

• THE FIELD IS y-DIRECTED, ONLY VARIES WITH z

$$\nabla \times \bar{\mathbf{H}} = \nabla \times \frac{\bar{\mathbf{B}}}{\mu_0} = -\frac{1}{\mu_0} \frac{\partial B_y}{\partial z} \hat{\mathbf{x}} = \frac{\partial \bar{\mathbf{D}}}{\partial t}$$

$$\bar{\mathbf{J}}_d = \frac{\partial \bar{\mathbf{D}}}{\partial t} = \frac{k_0 B_0}{\mu_0} \cos(\omega t) \cos(k_0 z) \hat{\mathbf{x}} \text{ A/m}^2$$

b. ELECTRIC FIELD INTENSITY

$$\bar{\mathbf{E}} = \frac{1}{\epsilon_0} \int \frac{\partial \bar{\mathbf{D}}}{\partial t} dt + C$$

$$= \frac{k_0 B_0}{\mu_0 \epsilon_0} \sin k_0 z \hat{\mathbf{x}} \int \cos \omega t dt$$

$$= \frac{k_0 B_0}{\omega \mu_0 \epsilon_0} \sin k_0 z \sin \omega t \hat{\mathbf{x}}$$

$C = 0$ because there is no steady electric field component when \mathbf{H} is time-varying

c. k_0

FIND THIS BY ENSURING CONSISTENCY BETWEEN THE TWO MAXWELL CURL EQUATIONS.

FARADAY'S LAW

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \Rightarrow \vec{B} = -\int \nabla \times \vec{E} dt$$

$C = 0$ because there is no steady magnetic field when E is time-varying.

FROM b), \vec{E} IS x -DIRECTED AND VARIES WITH z ONLY.

THUS, THE CURL OF \vec{E} SIMPLIFIES TO A SINGLE TERM.

$$\begin{aligned}\vec{B} &= -\int \nabla \times \vec{E} dt = -\int \frac{\partial E_z}{\partial z} \hat{y} dt \\ &= \int \frac{-k_0^2 B_0}{\omega \mu_0 \epsilon_0} \cos k_0 z \sin \omega t \hat{y} dt \\ &= \frac{k_0^2 B_0}{\omega^2 \mu_0 \epsilon_0} \cos k_0 z \cos \omega t \hat{y}\end{aligned}$$

$$\frac{k_0^2}{\omega^2 \mu_0 \epsilon_0} = 1 \Rightarrow k_0 = \omega \sqrt{\mu_0 \epsilon_0} \text{ m}^{-1}$$

We'll see this again in Chapter 11...

9.15 Use each of Maxwell's equations in point form to obtain as much information as possible about
 (a) \mathbf{H} , if $\mathbf{E} = 0$; (b) \mathbf{E} , if $\mathbf{H} = 0$.

THE TWO CASES CORRESPOND TO
 MAGNETOSTATIC AND ELECTROSTATIC
 SCENARIOS, RESPECTIVELY.

\mathbf{H} IF $\mathbf{E} = 0$

$$\nabla \times \mathbf{H} = \mathbf{J}_c$$

\mathbf{H} IS SOLELY DUE TO
 CONDUCTION CURRENT.

$$\frac{\partial \mathbf{B}}{\partial t} = 0$$

\mathbf{B} IS CONSTANT

$$\rho_v = 0$$

NO CHARGE IS
 PRESENT

$$\nabla \cdot \mathbf{B} = 0$$

ALWAYS TRUE

BECAUSE \mathbf{J} IS NON-ZERO BUT $\mathbf{E} = 0$,
 THE CURRENT CARRYING MEDIUM IS
 A PERFECT CONDUCTOR.

\mathbf{E} IF $\mathbf{H} = 0$

$$\nabla \times \mathbf{E} = 0$$

\mathbf{E} IS CONSTANT IN
 TIME & CONSERVATIVE

$$\nabla \times \mathbf{H} = \mathbf{J}_c + \frac{\partial \mathbf{D}}{\partial t} = 0 \Rightarrow \mathbf{J} = \frac{\partial \mathbf{D}}{\partial t} = 0$$

$$\nabla \cdot \mathbf{D} = \rho_v$$

\mathbf{E} ARISES FROM
 FREE CHARGE ONLY

$$\nabla \cdot \mathbf{B} = 0$$

$\mathbf{B} = 0$ ANYWAY.

9.18 The parallel-plate transmission line shown in Figure 9.7 has dimensions $b = 4$ cm and $d = 8$ mm, while the medium between the plates is characterized by $\mu_r = 1$, $\epsilon_r = 20$, and $\sigma = 0$. Neglect fields outside the dielectric. Given the field $\mathbf{H} = 5 \cos(10^9 t - \beta z) \mathbf{a}_y$ A/m, use Maxwell's equations to help find (a) β , if $\beta > 0$; (b) the displacement current density at $z = 0$; (c) the total displacement current crossing the surface $x = 0.5d$, $0 < y < b$, $0 < z < 0.1$ m in the \mathbf{a}_x direction.

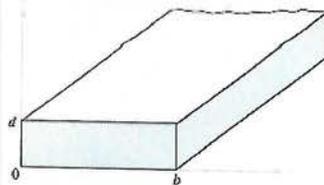


Fig. 9.7

$$\begin{aligned} \text{a. } \nabla \times \bar{\mathbf{H}} &= - \frac{\partial H_y}{\partial z} \hat{\mathbf{x}} = - 5\beta \sin(10^9 t - \beta z) \hat{\mathbf{x}} \\ &= 20\epsilon_0 \frac{\partial \bar{\mathbf{E}}}{\partial t} \end{aligned}$$

$$\begin{aligned} \bar{\mathbf{E}} &= \int \frac{-5\beta}{20\epsilon_0} \sin(10^9 t - \beta z) \hat{\mathbf{x}} dt \\ &= \frac{\beta}{4 \times 10^9 \epsilon_0} \cos(10^9 t - \beta z) \hat{\mathbf{x}} \end{aligned}$$

$$\begin{aligned} \nabla \times \bar{\mathbf{E}} &= \frac{\partial E_x}{\partial z} \hat{\mathbf{y}} = \frac{\beta^2}{4 \times 10^9 \epsilon_0} \sin(10^9 t - \beta z) \hat{\mathbf{y}} \\ &= -\mu_0 \frac{\partial \bar{\mathbf{H}}}{\partial t} \end{aligned}$$

$$\begin{aligned} \bar{\mathbf{H}} &= \int \frac{-\beta^2}{4 \times 10^9 \mu_0 \epsilon_0} \sin(10^9 t - \beta z) \hat{\mathbf{y}} dt \\ &= \frac{\beta^2}{4 \times 10^{18} \mu_0 \epsilon_0} \cos(10^9 t - \beta z) \hat{\mathbf{y}} \\ &= 5 \cos(10^9 t - \beta z) \hat{\mathbf{y}} \end{aligned}$$

$$\frac{\beta^2}{4 \times 10^{18} \mu_0 \epsilon_0} = 5$$

$$\beta^2 = \frac{4 \times 10^{18} \times 5}{(3 \times 10^8)^2}$$

$$\beta = \frac{\sqrt{2 \times 10^{19}}}{3 \times 10^8}$$

$$= 14.9 \text{ m}^{-1}$$

b. DISPLACEMENT CURRENT DENSITY @ $z=0$

$$\sigma = 0; \quad \nabla \times \bar{H} = \bar{J}_d = -5\beta \sin(10^9 t - \beta z) \hat{x}$$

$$= -74.5 \sin(10^9 t - \beta z) \hat{x}$$

$$= -74.5 \sin 10^9 t \hat{x} \frac{\text{A}}{\text{m}} \text{ at } z=0$$

c. TOTAL DISPLACEMENT CURRENT CROSSING THE SURFACE S

$$x = 0.5 \text{ d} = 4 \text{ mm} \quad 0 < y < b = 4 \text{ cm} \quad 0 < z < 0.1 \text{ m}$$

$$I_d = -74.5 b \int_0^{0.1} \sin(10^9 t - 14.9 z) \hat{x} \cdot \hat{z} dz$$

$$= \frac{-74.5 \times 0.04}{-14.9} \cos(10^9 t - 14.9 z) \Big|_0^{0.1}$$

$$= 0.20 [\cos(10^9 t - 1.49) - \cos(10^9 t)] \text{ A}$$

9.22

9.22 In a sourceless medium in which $\mathbf{J} = 0$ and $\rho_v = 0$, assume a rectangular coordinate system in which \mathbf{E} and \mathbf{H} are functions only of z and t . The medium has permittivity ϵ and permeability μ . (a) If $\mathbf{E} = E_x \mathbf{a}_x$ and $\mathbf{H} = H_y \mathbf{a}_y$, begin with Maxwell's equations and determine the second-order partial differential equation that E_x must satisfy. (b) Show that $E_x = E_0 \cos(\omega t - \beta z)$ is a solution of that equation for a particular value of β . (c) Find β as a function of given parameters.

a.

$$\nabla \times \bar{\mathbf{E}} = -\frac{\partial \bar{\mathbf{B}}}{\partial t} \Rightarrow \frac{\partial E_x}{\partial z} \hat{y} = -\mu \frac{\partial H_y}{\partial t} \hat{y}$$

$$\nabla \times \bar{\mathbf{H}} = -\frac{\partial \bar{\mathbf{D}}}{\partial t} \Rightarrow -\frac{\partial H_y}{\partial z} \hat{x} = \epsilon \frac{\partial E_x}{\partial t} \hat{x}$$

$$\frac{\partial^2 E_x}{\partial z^2} = -\mu \frac{\partial^2 H_y}{\partial t \partial z}$$

$$\frac{\partial^2 H_y}{\partial z \partial t} = -\epsilon \frac{\partial^2 E_x}{\partial t^2}$$

$$\frac{\partial^2 E_x}{\partial z^2} = \mu \epsilon \frac{\partial^2 E_x}{\partial t^2}$$

THE 2ND ORDER PDE
THAT E_x MUST
SATISFY.

b. SHOW THAT $E_x = E_0 \cos(\omega t - \beta z)$ IS
A SOLUTION

$$\frac{\partial^2 E_x}{\partial z^2} = -\beta^2 E_0 \cos(\omega t - \beta z)$$

$$\mu \epsilon \frac{\partial^2 E_x}{\partial t^2} = -\omega^2 \mu \epsilon E_0 \cos(\omega t - \beta z)$$

MUST
BE
EQUAL

c. $\beta^2 = \omega^2 \mu \epsilon$, $\beta = \omega \sqrt{\mu \epsilon}$

9.26 Write Maxwell's equations in point form in terms of \mathbf{E} and \mathbf{H} as they apply to a sourceless medium, where \mathbf{J} and ρ_v are both zero. Replace ϵ with μ , μ with ϵ , \mathbf{E} with \mathbf{H} , and \mathbf{H} with $-\mathbf{E}$, and show that the equations are unchanged. This is a more general expression of the *duality principle* in circuit theory.

$$\nabla \times \bar{\mathbf{E}} = -\mu \frac{\partial \bar{\mathbf{H}}}{\partial t} \quad (1)$$

$$\nabla \times \bar{\mathbf{H}} = \epsilon \frac{\partial \bar{\mathbf{E}}}{\partial t} \quad (2) \quad \text{MAXWELL'S EQUATIONS}$$

$$\nabla \cdot \epsilon \bar{\mathbf{E}} = 0 \quad (3) \quad \bullet \text{ POINT FORM}$$

$$\nabla \cdot \mu \bar{\mathbf{H}} = 0 \quad (4) \quad \bullet \text{ SOURCELESS MEDIUM}$$

$$\epsilon \rightarrow \mu \quad \mu \rightarrow \epsilon \quad \bar{\mathbf{E}} \rightarrow \bar{\mathbf{H}} \quad \bar{\mathbf{H}} \rightarrow -\bar{\mathbf{E}}$$

$$\nabla \times \bar{\mathbf{H}} = \epsilon \frac{\partial \bar{\mathbf{E}}}{\partial t} \quad (2)$$

$$\nabla \times \bar{\mathbf{E}} = -\mu \frac{\partial \bar{\mathbf{H}}}{\partial t} \quad (1)$$

$$\nabla \cdot \mu \bar{\mathbf{H}} = 0 \quad (4)$$

$$\nabla \cdot \epsilon \bar{\mathbf{E}} = 0 \quad (3)$$

$$(1) \rightarrow (2)$$

$$(2) \rightarrow (1)$$

$$(3) \rightarrow (4)$$

$$(4) \rightarrow (3)$$