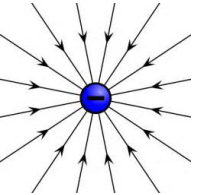
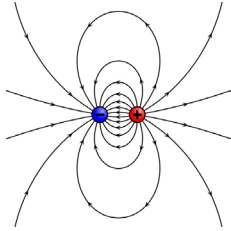


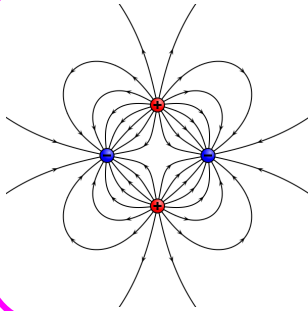
Lecture 11

$$V(\mathbf{r}) = \sum_{l=0}^{\infty} V_l(\mathbf{r}) = \boxed{V_0(\mathbf{r})} + \boxed{V_1(\mathbf{r})} + \boxed{V_2(\mathbf{r})} + \dots$$

monopole dipole quadrupole


$$V_0(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r}$$


$$V_1(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{\mathbf{p} \cdot \hat{\mathbf{r}}}{r^2}$$

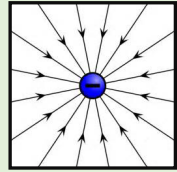

$$V_2(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \sum_{ij} \frac{Q_{ij} \hat{r}_i \hat{r}_j}{r^3}$$

Multipole expansion: Examples

Last Time: Multipole Expansion

- Monopole ($l = 0$)

$$V_0(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r}$$

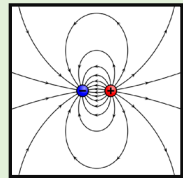


$$Q = \int_V \rho(\mathbf{r}') d\tau' \rightarrow \sum_a q_a$$

$\mathbf{E}_0(\mathbf{r})$

- Dipole ($l = 1$)

$$V_1(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{\mathbf{p} \cdot \hat{\mathbf{r}}}{r^2}$$

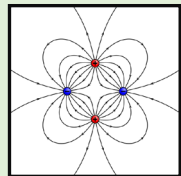


$\mathbf{E}_1(\mathbf{r})$

$$\mathbf{p} = \int_V \rho(\mathbf{r}') \mathbf{r}' d\tau' \rightarrow \sum_a q_a \mathbf{r}_a$$

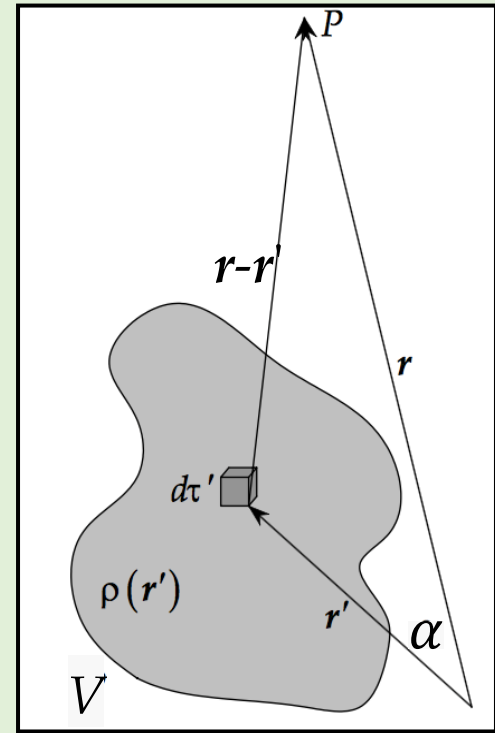
- Quadrupole ($l = 2$)

$$V_2(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \sum_{ij} \frac{Q_{ij} \hat{r}_i \hat{r}_j}{r^3}$$



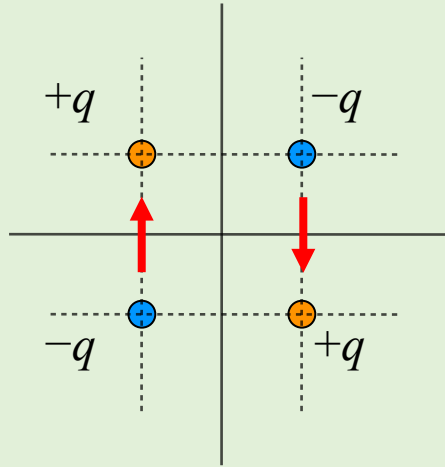
$\mathbf{E}_2(\mathbf{r})$

$$Q_{ij} = \int_V \frac{\rho(\mathbf{r}')}{2} (3r'_i r'_j - r'^2 \delta_{ij}) d\tau' \rightarrow \sum_a \frac{q_a}{2} (3r_{a,i} r_{a,j} - r_a^2 \delta_{ij})$$



Last Time: Multipole Expansion

Q: What can you say about the potential of the following charge distribution?



$$Q = \sum_a q_a = 0 \rightarrow V_0(\mathbf{r}) = 0$$

$$\mathbf{p} = \sum_a q_a \mathbf{r}_a = 0 \rightarrow V_1(\mathbf{r}) = 0$$

A. $V(\mathbf{r}) = V_0(\mathbf{r}) + \text{higher-order terms}$

B. $V(\mathbf{r}) = V_1(\mathbf{r}) + \text{higher-order terms}$

☒ C. $V(\mathbf{r}) = V_2(\mathbf{r}) + \text{higher-order terms}$

D. $V(\mathbf{r}) = 0$

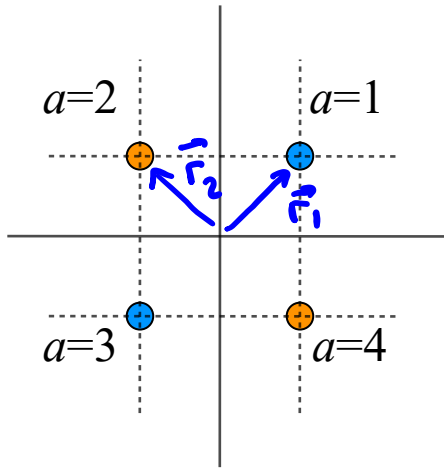
E. None of the above.

$$Q_{ij} = \sum_a \frac{q_a}{2} (3r_{a,i}r_{a,j} - r_a^2 \delta_{ij})$$

$$V_2(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \sum_{ij} \frac{Q_{ij} \hat{r}_i \hat{r}_j}{r^3}$$

Quadrupole Moment & Potential

Q: Compute the quadrupole moment, Q_{ij} . Say we have point charges with $(x, y) = (\pm d, \pm d)$:



$$Q_{ij} = \int_V \frac{\rho(\mathbf{r}')}{2} (3r'_i r'_j - r'^2 \delta_{ij}) d\tau' \rightarrow \sum_a \frac{q_a}{2} (3r_{a,x} r_{a,y} - r_a^2 \delta_{xy})$$

$$\mathbf{Q} = \begin{pmatrix} Q_{xx} & Q_{xy} & Q_{xz} \\ Q_{yx} & Q_{yy} & Q_{yz} \\ Q_{zx} & Q_{zy} & Q_{zz} \end{pmatrix}$$

a = index counting charges; $a = 1, 2, 3, 4$

\mathbf{r}_a = position of a -th charge

$r_{a,i}$ = i -th coordinate of a -th charge, $i = x, y, z$

$$\delta_{ij} = \begin{cases} 1, & i=j \\ 0, & i \neq j \end{cases}$$

$$\delta_{xx} = 1$$

$$\delta_{xy} = 0$$

Kronecker's delta

Q_{xx} :

A. Got the value

$$i=x \quad j=x$$

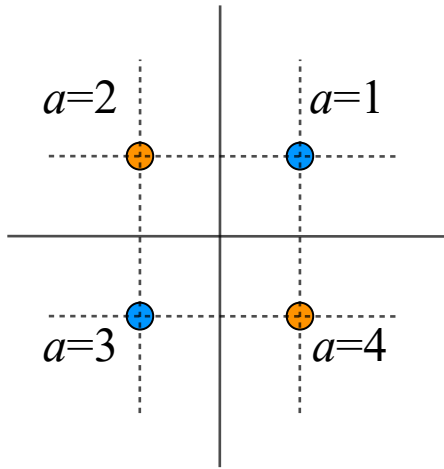
B. Set the sum up

C. ☹

$$Q_{yy} \rightarrow Q_{zz}$$

Quadrupole Moment & Potential

Q: Compute the quadrupole moment, Q_{ij} . Say we have point charges with $(x, y) = (\pm d, \pm d)$:



$$Q_{ij} = \int_V \frac{\rho(\mathbf{r}')}{2} (3r'_i r'_j - r'^2 \delta_{ij}) d\tau' \rightarrow \sum_a \frac{q_a}{2} (3r_{a,i} r_{a,j} - \cancel{r_a^2 \delta_{ij}})$$

$3 \cdot x_a \cdot y_a$ $\delta_{xy} \equiv 0$

$$r_a^2 = x_a^2 + y_a^2 + z_a^2 = 2d^2$$

$$(x_a = y_a = d, \quad z_a = 0)$$

$$\bullet Q_{xx} = \sum_a \frac{q_a}{2} (3x_a^2 - r_a^2) = \frac{d^2}{2} \sum_a q_a = 0$$

$$\bullet \text{ Similarly, } Q_{yy} = 0$$

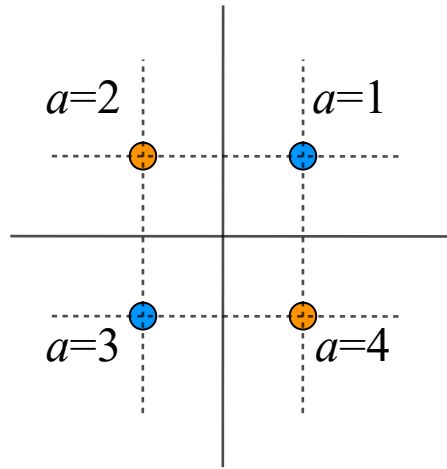
$$\bullet Q_{zz} = \sum_a \frac{q_a}{2} (0 - r_a^2) = -2d^2 \sum_a \frac{q_a}{2} = 0$$

- All diagonal elements of the quadrupole moment are equal to zero!

✓ Q_{xy} Q_{yz} Q_{zx} Q_{yx} Q_{zy} Q_{xz}

Quadrupole Moment & Potential

Q: Compute the quadrupole moment, Q_{ij} . Say we have point charges with $(x, y) = (\pm d, \pm d)$:



$$Q_{ij} = \int_V \frac{\rho(\mathbf{r}')}{2} (3r'_i r'_j - r'^2 \delta_{ij}) d\tau' \rightarrow \sum_a \frac{q_a}{2} (3r_{a,i} r_{a,j} - r_a^2 \delta_{ij})$$

$$Q_{xz}: 3x_a z_a - r_a^2 \delta_{xz}$$

$$\begin{aligned} \bullet Q_{xy} &= \sum_{a=1}^4 \frac{q_a}{2} 3x_a y_a = Q_{yx} = \\ &= \sum_a \frac{3}{2} [(-q)(d)(d) + (q)(-d)(d) + (-q)(-d)(-d) + (q)(d)(-d)] \end{aligned}$$

$$\frac{1}{4\pi\epsilon_0} \left\{ \frac{Q_{xy} xy}{r^5} + \frac{Q_{yx} yx}{r^5} \right\} = -6qd^2$$

$$i, j = x, y, z$$

• Potential:

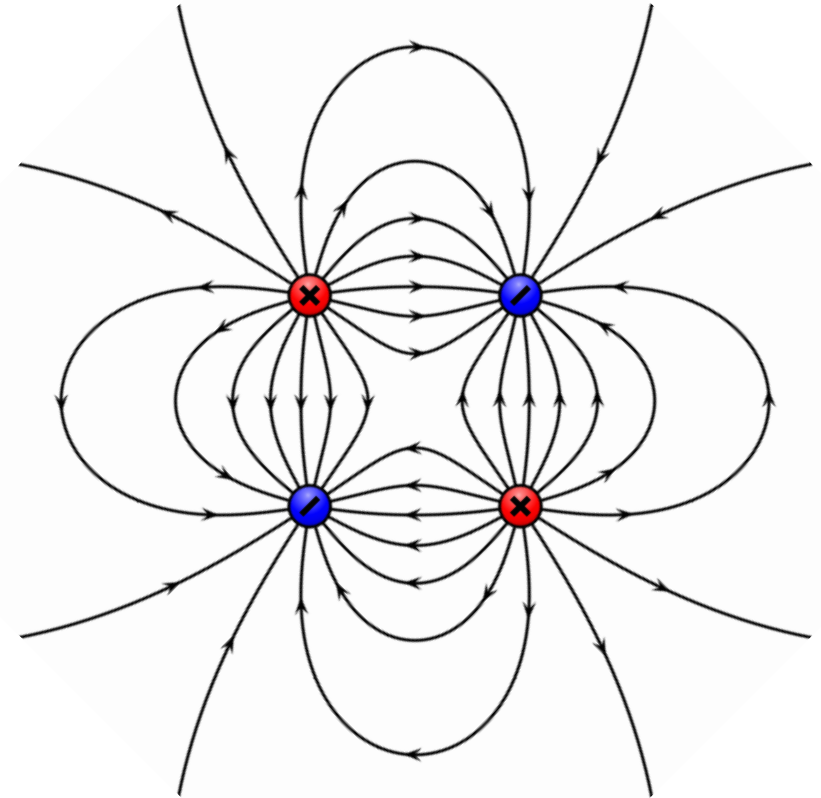
$$\mathbf{Q} = \begin{pmatrix} 0 & -6 & 0 \\ -6 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} qd^2 \quad (z_a = 0)$$

$$\rightarrow V_2(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \sum_{ij} \frac{Q_{ij} r_i r_j}{r^5} = -\frac{12qd^2}{4\pi\epsilon_0} \frac{xy}{r^5} \propto \frac{1}{r^3} \text{ at large } r$$

Quadrupole Field Geometry

The electric field is readily obtained by taking the gradient of $V(\mathbf{r})$:

$$\mathbf{E}_l(\mathbf{r}) = -\nabla V_l(\mathbf{r}) \quad (l = 0, 1, 2, \dots)$$



Exercise 1: Multipoles of a Charged Ring

A circular ring of charge centered at the origin with radius R carries a linear charge density λ . Find the first three terms ($l = 0, 1, 2$) in the multipole expansion for $V(\mathbf{r})$.

Reminder:

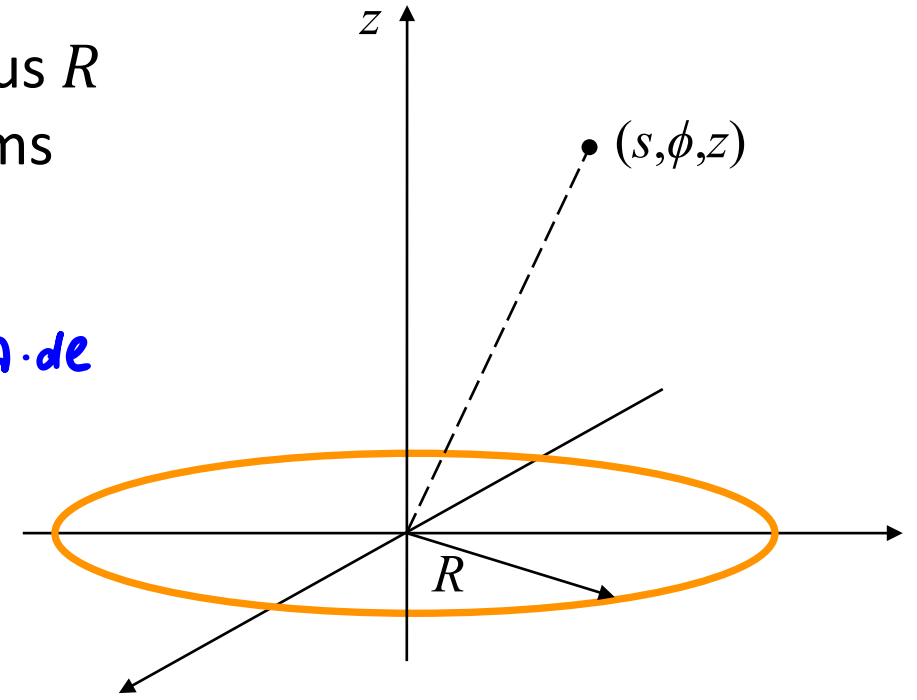
$$Q = Q_{\text{tot}} = \lambda \cdot 2\pi R$$

$$\rho(\mathbf{r}') d\tau' = \lambda \cdot dl$$

A. $Q \equiv \int_V \rho(\mathbf{r}') d\tau'$ $V_0(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r}$

B. $\mathbf{p} \equiv \int_V \rho(\mathbf{r}') \mathbf{r}' d\tau'$ $V_1(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{\mathbf{p} \cdot \hat{\mathbf{r}}}{r^2}$

C. $Q_{ij} \equiv \int_V \frac{\rho(\mathbf{r}')}{2} (3r'_i r'_j - r'^2 \delta_{ij}) d\tau'$ $V_2(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \sum_{ij} \frac{Q_{ij} \hat{r}_i \hat{r}_j}{r^3}$



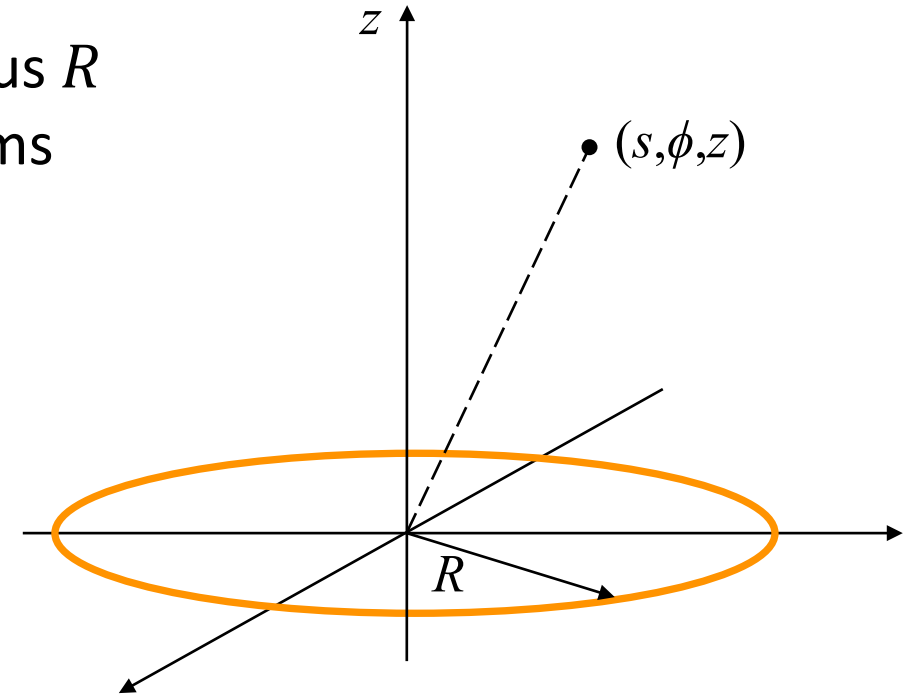
Exercise 1: Multipoles of a Charged Ring

A circular ring of charge centered at the origin with radius R carries a linear charge density λ . Find the first three terms ($l = 0, 1, 2$) in the multipole expansion for $V(\mathbf{r})$.

Monopole ($l = 0$):

$$Q = \int_V \rho(\mathbf{r}') d\tau'$$

$$\rightarrow \int_C \lambda(\mathbf{r}') dl' = \lambda \int_0^{2\pi} R d\phi' = 2\pi R \lambda$$



$$V_0(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r} = \frac{\lambda}{2\epsilon_0} \frac{R}{\sqrt{s^2 + z^2}}$$

Exercise 1: Multipoles of a Charged Ring

A circular ring of charge centered at the origin with radius R carries a linear charge density λ . Find the first three terms ($l = 0, 1, 2$) in the multipole expansion for $V(\mathbf{r})$.

Dipole ($l = 1$): $\mathbf{p} \equiv \int_V \rho(\mathbf{r}') \mathbf{r}' d\tau'$

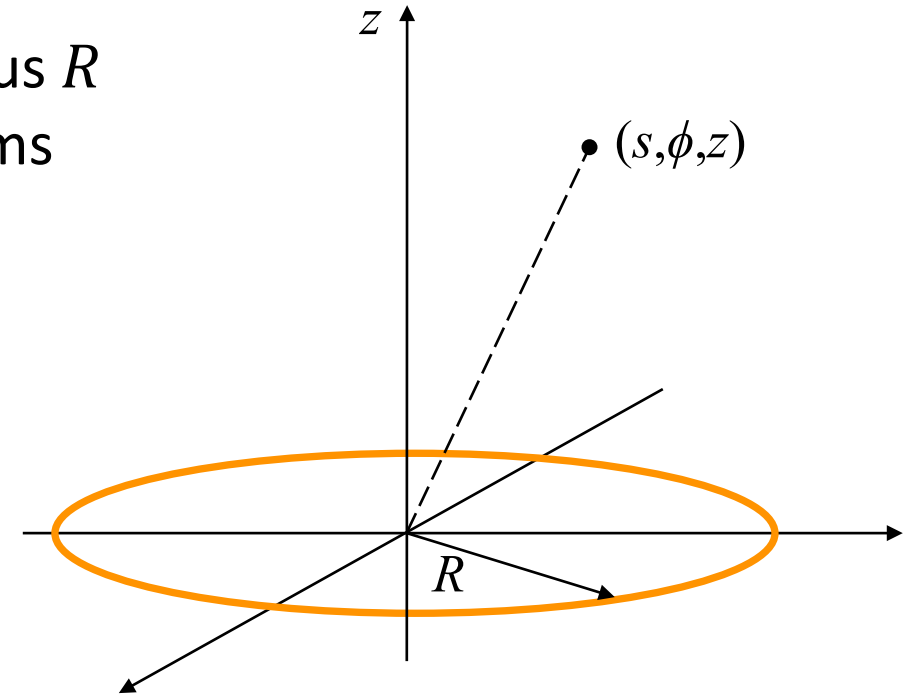
Q: Which coordinate system do you prefer to express \mathbf{p} ?

- Cartesian coordinate system
=> you will get (p_x, p_y, p_z)

$$\mathbf{r}' = x' \hat{\mathbf{x}} + y' \hat{\mathbf{y}} + z' \hat{\mathbf{z}}$$

- Cylindrical coordinate system
=> you will get (p_s, p_z)

$$\mathbf{r}' = s' \hat{\mathbf{s}} + z' \hat{\mathbf{z}}$$



- A. Zero
- B. $2\pi R^2 \lambda \hat{\mathbf{s}}$
- C. Something else
- D. ☹️

Q: Pick a coordinate system, compute \mathbf{p} and submit your answer

Exercise 1: Multipoles of a Charged Ring

Dipole ($l = 1$): $\mathbf{p} \equiv \int_V \rho(\mathbf{r}') \mathbf{r}' d\tau'$ $\rho(r')d\tau' = \lambda dl = \lambda R d\phi$

• Cartesian: $\mathbf{r}' = x'\hat{\mathbf{x}} + y'\hat{\mathbf{y}} + z'\hat{\mathbf{z}}$

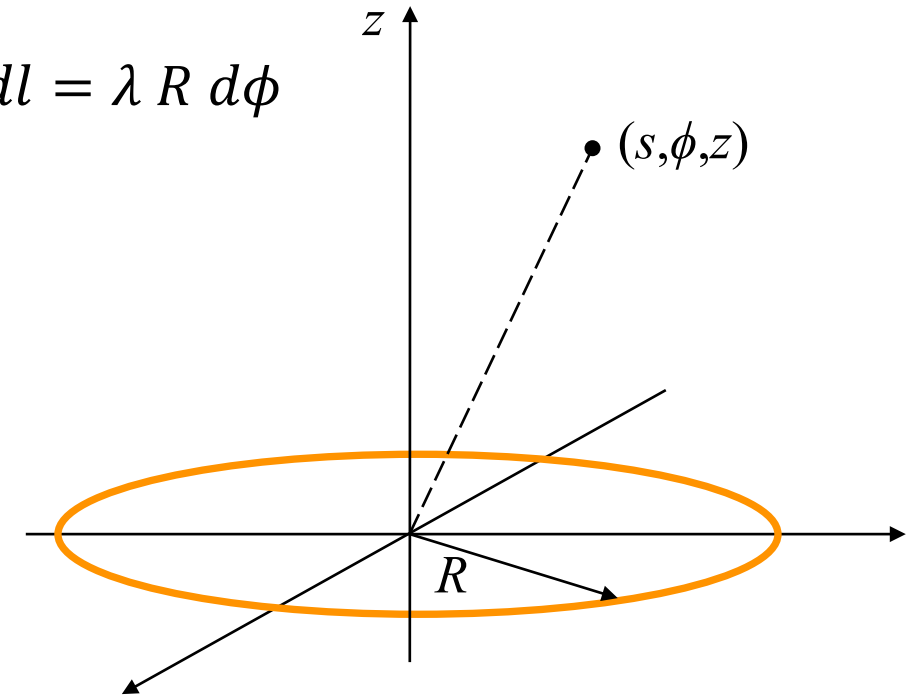
$$p_x = \int_C \lambda(\mathbf{r}') x' dl' = \lambda \int_0^{2\pi} R \cos \phi' R d\phi' = 0$$

Similarly for p_y . And $p_z \equiv 0$ since $z' \equiv 0$.

• Cylindrical: $\mathbf{r}' = s'\hat{\mathbf{s}} + z'\hat{\mathbf{z}}$

$$p_s = \int_C \lambda(\mathbf{r}') s' dl' = \lambda \int_0^{2\pi} R \cdot R d\phi' = \lambda \cdot 2\pi R^2$$

...and $p_z \equiv 0$ since $z' \equiv 0$.



- A. Zero
- B. $2\pi R^2 \lambda \hat{\mathbf{s}}$
- C. Something else
- D. ☹️

Q1: What's going on ???

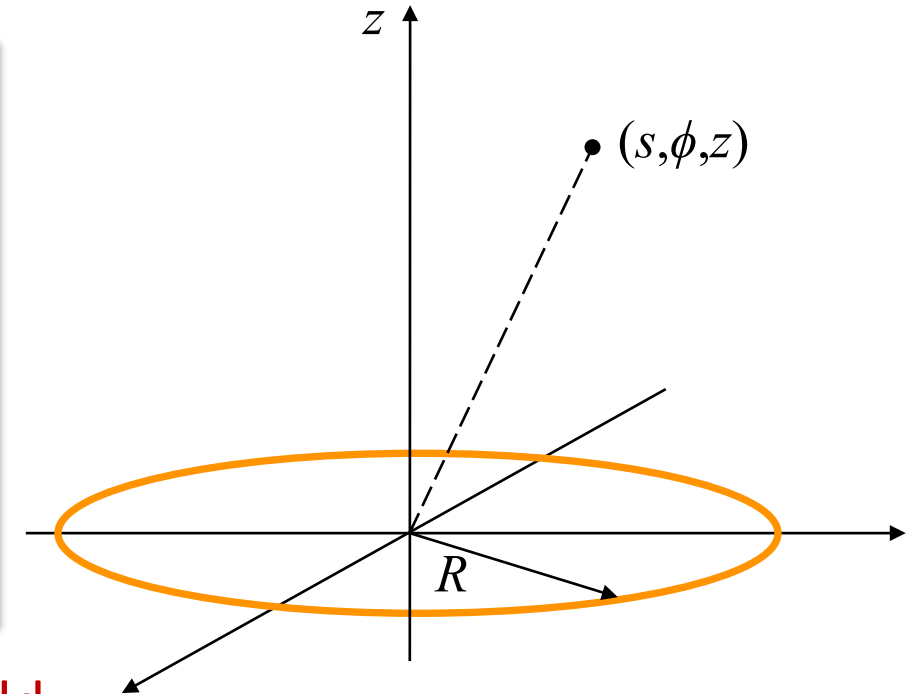
Q2: Which answer you expect is true?

Exercise 1: Multipoles of a Charged Ring

- We need to recall what we talked about at Lecture 1:

Warning

- Cartesian coordinate system:
 $\mathbf{r}_1 = (x_1, y_1, z_1)$ $r_1 \equiv |\mathbf{r}_1| = \sqrt{x_1^2 + y_1^2 + z_1^2}$
 $\mathbf{r}_2 = (x_2, y_2, z_2)$ $\mathbf{R} = \mathbf{r}_1 + \mathbf{r}_2 = (x_1 + x_2, y_1 + y_2, z_1 + z_2)$
- Cylindrical coordinate system:
 $\mathbf{r}_1 = (s_1, \theta_1, z_1)$ ~~$r_1 \equiv |\mathbf{r}_1| \neq \sqrt{s_1^2 + \theta_1^2 + z_1^2}$~~
 $\mathbf{r}_2 = (s_2, \theta_2, z_2)$ ~~$\mathbf{R} = \mathbf{r}_1 + \mathbf{r}_2 \neq (s_1 + s_2, \theta_1 + \theta_2, z_1 + z_2)$~~
- Spherical coordinate system: it does **not** work the “Cartesian” way either!



- Integration is a form of vector addition. We only can add components of Cartesian vectors, to find their sum (since this coordinate system has fixed unit vectors).
- You can carry integration out in any coordinate system, but you need to set up integrals for Cartesian component of the vector you are looking for.



A.

Zero

B.

$2\pi R^2 \lambda \hat{s}$

C. Something else

D. ☹️

Exercise 1: Multipoles of a Charged Ring

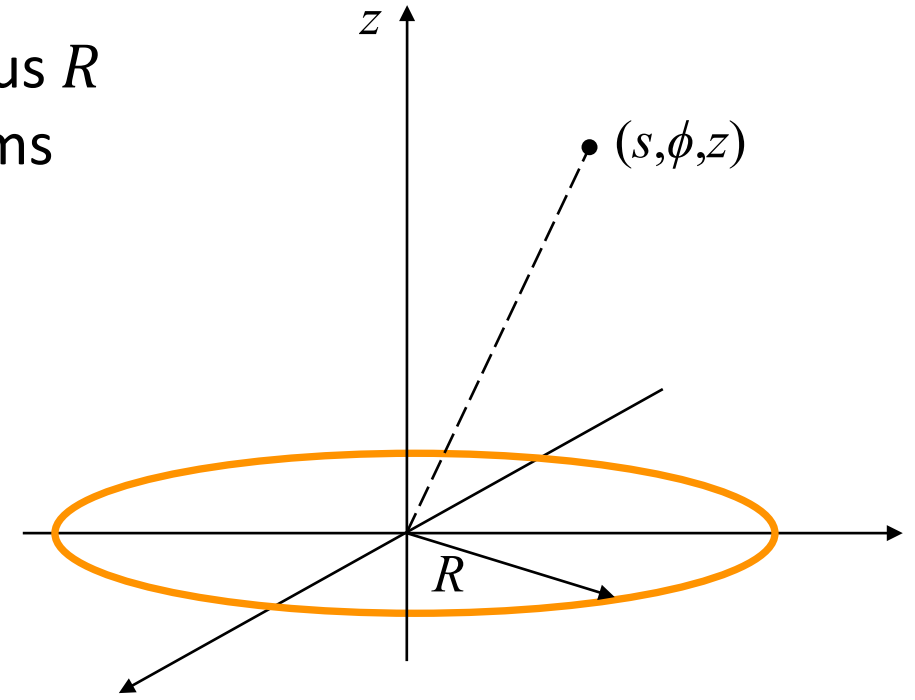
A circular ring of charge centered at the origin with radius R carries a linear charge density λ . Find the first three terms ($l = 0, 1, 2$) in the multipole expansion for $V(\mathbf{r})$.

Quadrupole ($l = 2$):

$$Q_{ij} = \int_C \frac{\lambda(\mathbf{r}')}{2} (3 \overset{x' x'}{r'_i r'_j} - \overset{R^2 \delta_{xx}}{r'^2 \delta_{ij}}) dl' \quad dl' = R d\phi'$$

In the ring: $r'_x = x' = R \cos \phi'$, $r'_y = y' = R \sin \phi'$
 $r'_z = z' = 0$, $r'^2 = x'^2 + y'^2 + z'^2 = R^2$

$$Q_{xx} = \frac{\lambda R^3}{2} \int_0^{2\pi} (3 \cos^2 \phi' - 1) d\phi' = \boxed{\frac{\pi \lambda R^3}{2}}$$



$Q_{yy} = Q_{xx}$ by symmetry

- Exercise:
Work out Q_{zz}

Exercise 1: Multipoles of a Charged Ring

A circular ring of charge centered at the origin with radius R carries a linear charge density λ . Find the first three terms ($l = 0, 1, 2$) in the multipole expansion for $V(\mathbf{r})$.

Quadrupole ($l = 2$):

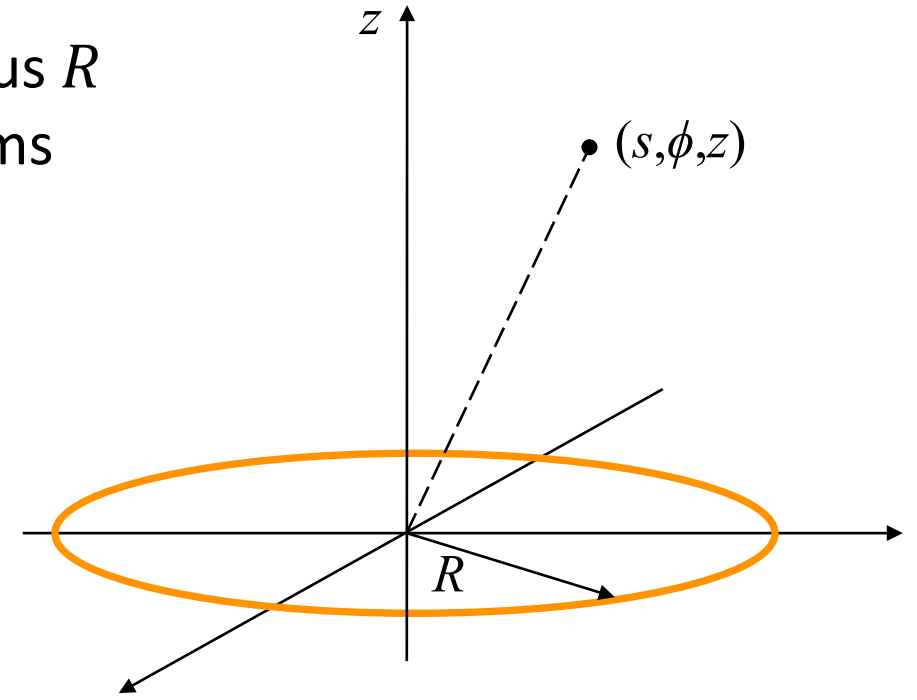
$$Q_{zz} = \frac{\lambda R^3}{2} \int_0^{2\pi} (3 \cdot 0 - 1) d\phi' = \boxed{-\pi \lambda R^3}$$

$$Q_{xy} = Q_{xz} = Q_{yz} = 0 \quad (\text{left as an exercise})$$

$$\rightarrow V_2(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \left(\frac{Q_{xx}x^2}{r^5} + \frac{Q_{yy}y^2}{r^5} + \frac{Q_{zz}z^2}{r^5} \right)$$

$$= \frac{\pi \lambda R^3}{8\pi\epsilon_0} \frac{x^2 + y^2 - 2z^2}{r^5}$$

$$= \frac{Q}{4\pi\epsilon_0} \frac{R^2}{4} \frac{s^2 - 2z^2}{(s^2 + z^2)^{5/2}} \quad (Q = 2\pi R\lambda)$$



- Now let's see what this looks like in the x - y plane, and along the z axis.

$$V_0(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r} = \frac{\lambda}{2\epsilon_0} \frac{R}{\sqrt{s^2 + z^2}}$$

Exercise 1: Multipoles of a Charged Ring

$$V_2(\mathbf{r}) = \frac{Q}{4\pi\epsilon_0} \frac{R^2}{4} \frac{s^2 - 2z^2}{(s^2 + z^2)^{5/2}}$$

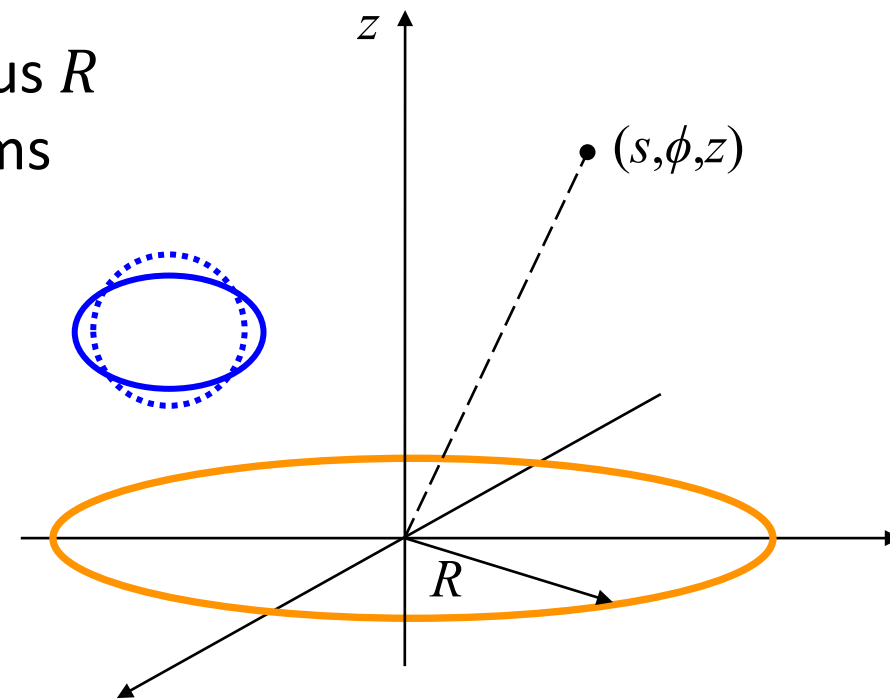
A circular ring of charge centered at the origin with radius R carries a linear charge density λ . Find the first three terms ($l = 0, 1, 2$) in the multipole expansion for $V(\mathbf{r})$.

- In the x - y plane:
 $z = 0$ and $s = r$:

$$V(\mathbf{r}) \rightarrow \frac{Q}{4\pi\epsilon_0} \left(\frac{1}{s} + \frac{R^2}{4s^3} \right) > V_0(\mathbf{r})$$

- Along the z axis:
 $s = 0$ and $|z| = r$:

$$V(\mathbf{r}) \rightarrow \frac{Q}{4\pi\epsilon_0} \left(\frac{1}{|z|} - \frac{R^2}{2|z|^3} \right) < V_0(\mathbf{r})$$



- In the x - y plane, the quadrupole term gives an **increment** to the monopole potential at a given r , while along the z axis it gives a **decrement** to the monopole potential.
- This is because this charge distribution differs from a “point-like” charge at the origin by “spilling” the charge density out in the x - y plane and away from the z -axis.

A Slightly Different Approach

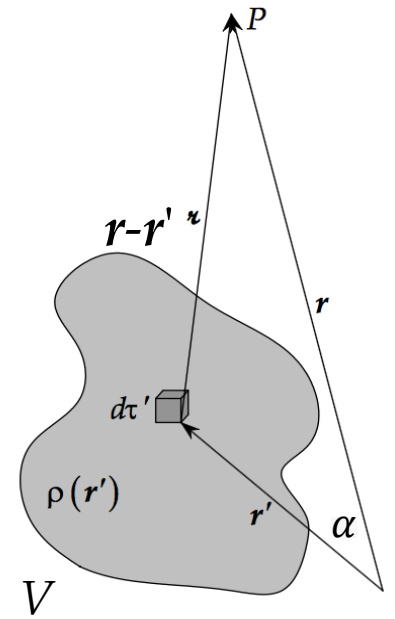
- Recall: We started from here...

We get:

$$\frac{1}{|\mathbf{r} - \mathbf{r}'|} = \frac{1}{r} \left[1 + \frac{r'}{r} \cos\alpha + \frac{1}{2} \frac{r'^2}{r^2} (3 \cos^2\alpha - 1) + \mathcal{O}\left(\frac{r'^3}{r^3}\right) \right]$$

Now we can expand Coulomb's law in powers of $1/r^{l+1}$:

$$V(\mathbf{r}) = \sum_{l=0}^{\infty} V_l(\mathbf{r}) = \underbrace{V_0(\mathbf{r})}_{\text{monopole}} + \underbrace{V_1(\mathbf{r})}_{\text{dipole}} + \underbrace{V_2(\mathbf{r})}_{\text{quadrupole}} + \dots \underbrace{\dots}_{\text{octupole+}}$$



- ...and split $V_l(\mathbf{r})$ into two parts, one depends only on \mathbf{r}' (moments) and the other only on \mathbf{r} :

$$Q \equiv \int_V \rho(\mathbf{r}') d\tau'$$

$$\mathbf{p} \equiv \int_V \rho(\mathbf{r}') \mathbf{r}' d\tau'$$

$$Q_{ij} \equiv \int_V \frac{\rho(\mathbf{r}')}{2} (3r'_i r'_j - r'^2 \delta_{ij}) d\tau'$$

$$V_0(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{Q}{r}$$

$$V_1(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{\mathbf{p} \cdot \hat{\mathbf{r}}}{r^2}$$

$$V_2(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \sum_{ij} \frac{Q_{ij} \hat{r}_i \hat{r}_j}{r^3}$$

A Slightly Different Approach

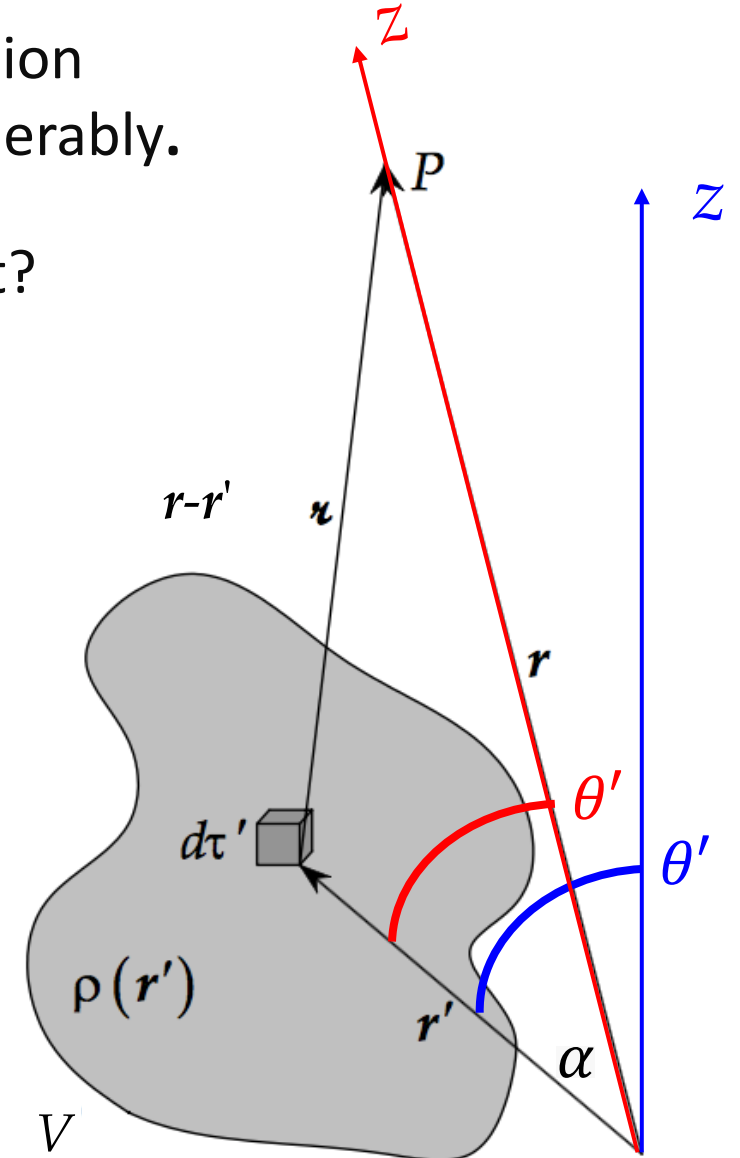
- We did this to eliminate the angle α , which is NOT the integration angle, θ' . This angle complicates the integration over \mathbf{r}' considerably.
- But, what if we rotate the z-axis towards the observation point?
- Now $\alpha = \theta'$, and our original expansion can be used directly:

$$\frac{1}{|\mathbf{r} - \mathbf{r}'|} = \frac{1}{r} \left[1 + \frac{r'}{r} \cos \theta' + \frac{1}{2} \frac{r'^2}{r^2} (3 \cos^2 \theta' - 1) + \mathcal{O} \left(\frac{r'^3}{r^3} \right) \right]$$

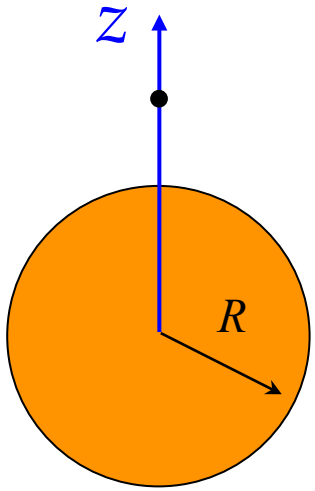
$$V(\mathbf{r}) = \sum_{l=0}^{\infty} V_l(\mathbf{r}) = \underbrace{V_0(\mathbf{r})}_{\text{monopole}} + \underbrace{V_1(\mathbf{r})}_{\text{dipole}} + \underbrace{V_2(\mathbf{r})}_{\text{quadrupole}} + \dots \underbrace{\dots}_{\text{octupole+}}$$

$$V(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int_V \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\tau'$$

Don't need to worry about, e.g., which components of \mathbf{p} to take, just compute these integrals!



Exercise 2: Multipoles of a Charged Sphere



A uniformly charged sphere carries a volume charge density ρ . Find the first three terms ($l = 0, 1, 2$) in the multipole expansion for $V(\mathbf{r})$.

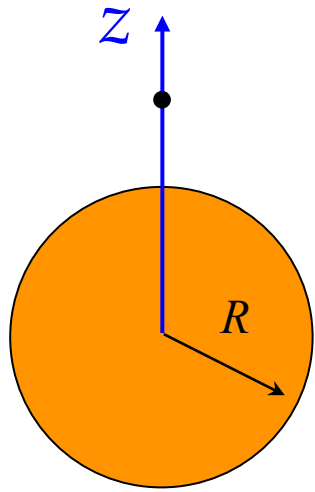
$$\frac{4\pi}{3} R^3 \rho \quad \swarrow \text{uniform!}$$

- Due to the central symmetry of the charge distribution, we can assume that our observation point is on the z -axis – the answer that we will get would be the same for all other directions

$$V_0(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{1}{r} \left(\int_V \rho(\mathbf{r}') d\tau' \right)$$

$$V_1(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{1}{r^2} \int_V \rho(\mathbf{r}') r' \cos \theta' d\tau'$$

$$V_2(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{1}{r^3} \int_V \rho(\mathbf{r}') \frac{r'^2}{2} (3 \cos^2 \theta' - 1) d\tau'$$



Exercise 2: Multipoles of a Charged Sphere

A uniformly charged sphere carries a volume charge density ρ . Find the first three terms ($l = 0, 1, 2$) in the multipole expansion for $V(\mathbf{r})$.

$$\bullet V_0(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{1}{r} \int_V \rho(\mathbf{r}') d\tau' = \frac{1}{4\pi\epsilon_0} \frac{Q}{r} \quad Q = \rho \frac{4\pi}{3} R^3$$

$$\bullet V_1(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{1}{r^2} \int_V \rho(\mathbf{r}') r' \cos \theta' d\tau' = \frac{1}{4\pi\epsilon_0} \frac{1}{r^2} \int_V (\rho r' \cos \theta') (r'^2 \sin \theta' dr' d\theta' d\phi') = 0$$

$$\bullet V_2(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{1}{r^3} \int_V \rho(\mathbf{r}') \frac{r'^2}{2} (3 \cos^2 \theta' - 1) d\tau'$$

$$= \frac{1}{4\pi\epsilon_0} \frac{1}{r^3} \int_V \left(\frac{\rho}{2} \frac{r'^2}{2} (3 \cos^2 \theta' - 1) \right) (r'^2 \sin \theta' dr' d\theta' d\phi') = 0$$

Why ???

Choice of Origin in Multipole Expansion

- Consider a point charge at the origin:

$$V(r) = \frac{kq}{r} \quad (\text{Monopole only, of course})$$

...and one not at the origin, say at $(r, \theta, \varphi) = (a, \theta_0, 0)$:

$$V(r) = \frac{kq}{d} = \frac{kq}{|\vec{r} - \vec{r}_a|}$$

$$V = \frac{q}{r} \left[1 + \frac{a}{r} \cos \theta_0 + \left(\frac{a}{r} \right)^2 P_2(\cos \theta_0) + \left(\frac{a}{r} \right)^3 P_3(\cos \theta_0) + \dots \right]$$

Q: What is going on here? Why does a point charge have so many non-zero moments all of a sudden?

- Changing the origin doesn't change the physics (and you will show it in HW-3!), but it can radically change the terms in the series.
- It's really important to specify origin when calculating moments!
- Dipole moment independent of origin only if $Q = 0$!