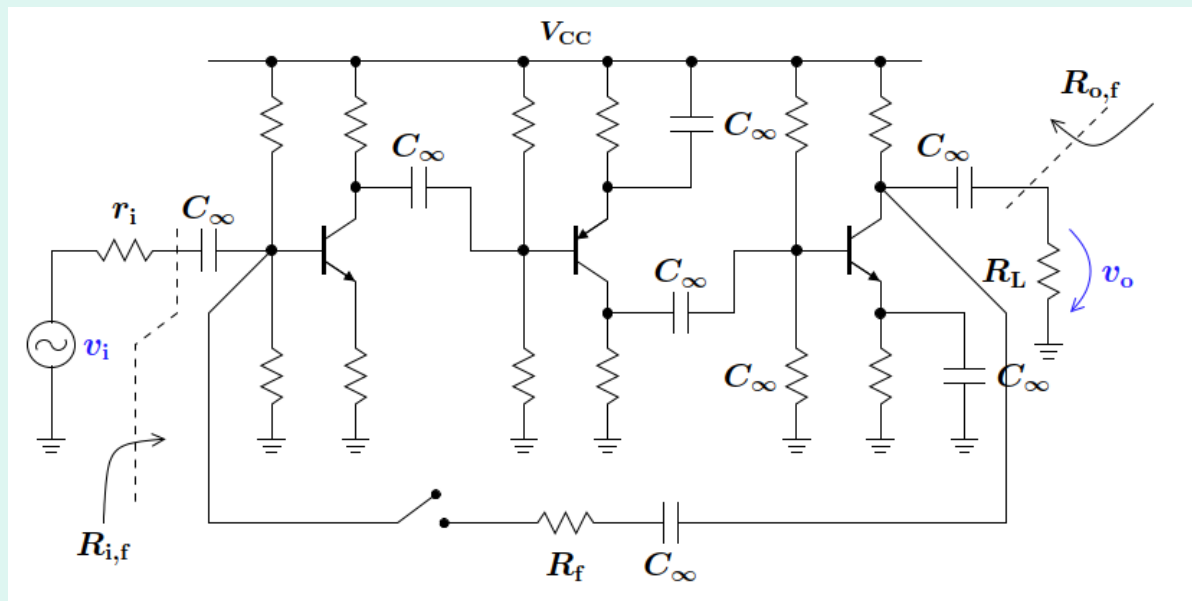
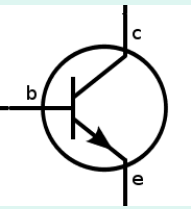


ELEC 301 - Band-pass circuits

L08 - Sep 22

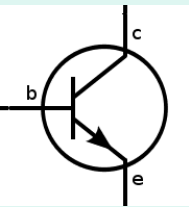
Instructor: Edmond Cretu





Last time

- Bode plots - log magnitude and phase plots
- Examples
- Approximation errors



L08 Q01 - Bode plots for a 2nd order system

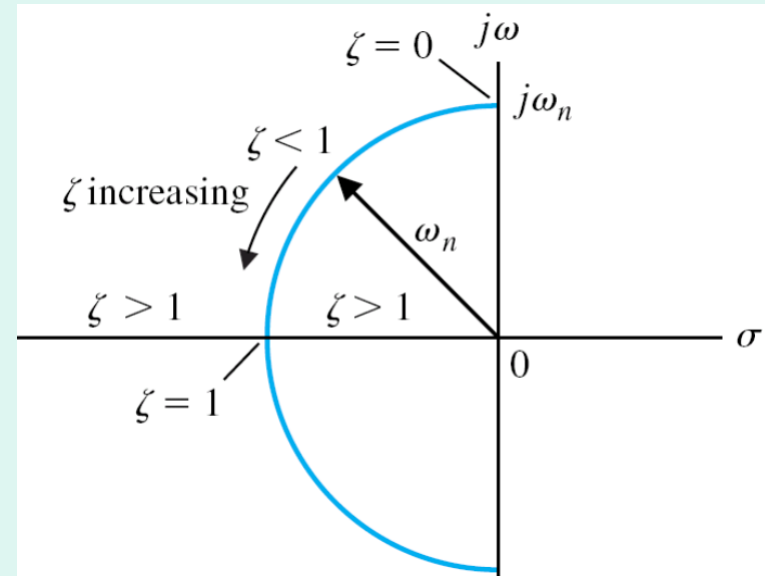
- Given a second order system, in which case the Bode technique provides a better approximation of the real frequency response?

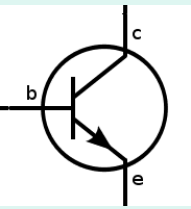
A. quality factor $Q=1000$

B. quality factor $Q=10$

C. Quality factor $Q=0.5$

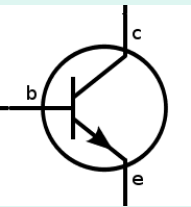
$$Y(s) = \frac{\omega_n^2}{s^2 + 2\zeta\omega_n s + \omega_n^2} X(s)$$





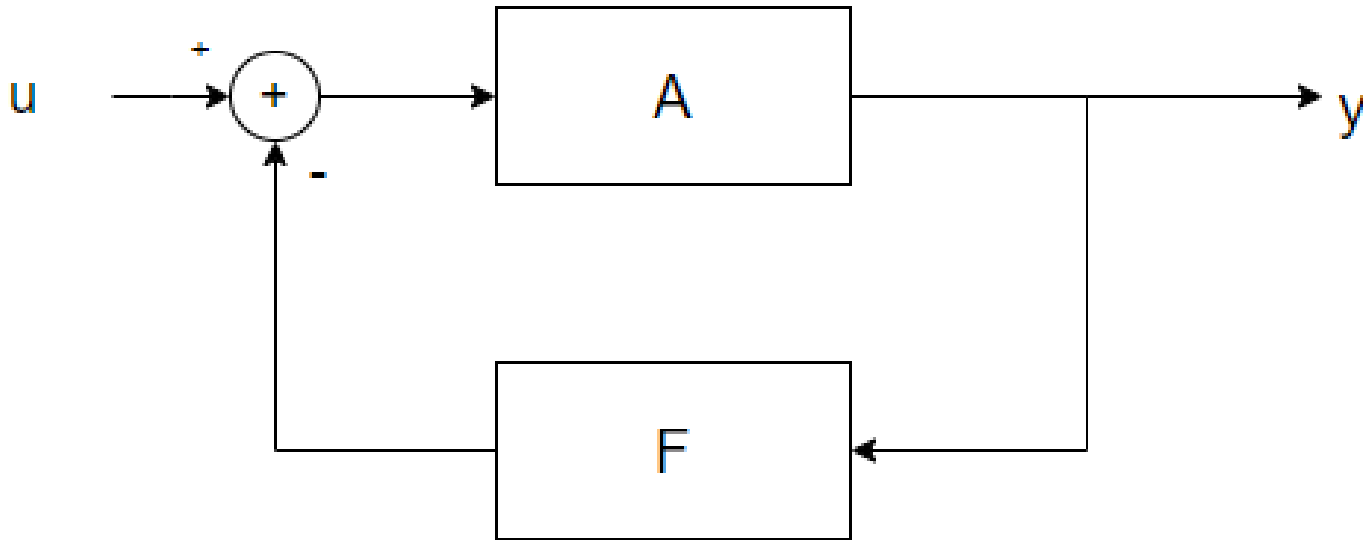
Conclusions

- Bode plots is a fast visual approximation technique for the frequency response of a linear system
- It provides a good approximation for systems with distinct individual poles on the real axis
- The typical cases for electronic amplifiers correspond to such a situation

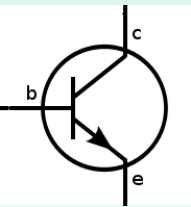


Simple LP active amplifier

- Simple low-pass amplifier stage with passive feedback network

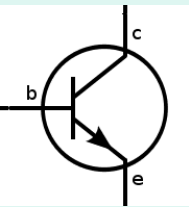


$$A(j\omega) = \frac{A_0}{1 + j \frac{\omega}{\omega_H}}, \quad F = F_0 = \frac{R_1}{R_1 + R_2} < 1$$



L08 Q02 LP amplifier

- Which statement is **false** for the previous configuration?
 - A. The closed loop gain decreases because of the negative feedback
 - B. The critical frequency of the closed loop transfer function decreases as result of the negative feedback action
 - C. The critical frequency of the closed loop transfer function increases, compared to the open loop case
 - D. The closed loop system is less sensitive to errors in the value of A_0

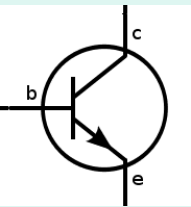


Negative feedback action

- The critical frequency ω_H is increased by the loop gain

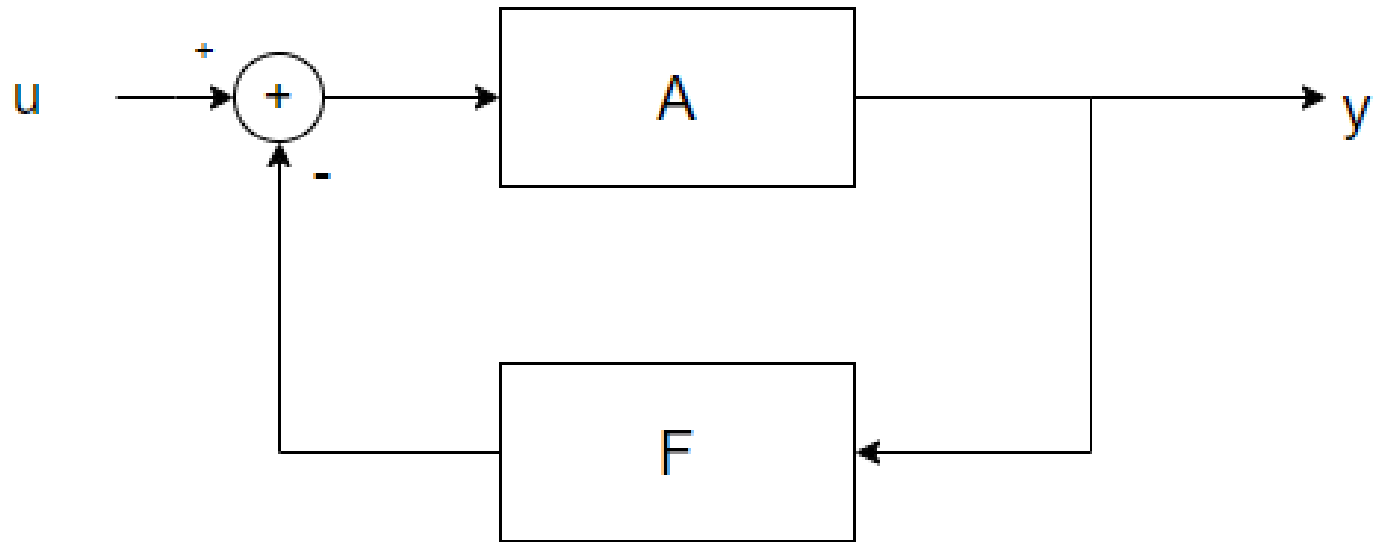
$$A_f(j\omega) = \frac{A(j\omega)}{1 + A(j\omega)F(j\omega)} = \frac{\frac{A_0}{1 + j\frac{\omega}{\omega_H}}}{1 + \frac{A_0}{1 + j\frac{\omega}{\omega_H}}F_0} = \frac{A_0}{1 + j\frac{\omega}{\omega_H} + A_0F_0}$$

$$A_f(j\omega) = \frac{A_0}{1 + A_0F_0} \frac{1}{1 + j\frac{\omega}{\omega_H(1 + A_0F_0)}} \stackrel{A_0F_0 \gg 1}{\approx} \frac{1}{F_0} \frac{1}{1 + j\frac{\omega}{\omega_H(1 + A_0F_0)}}$$

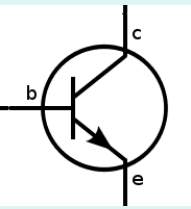


High-pass amplifier case

- Similar active network, but with a high-pass amplifier characteristics

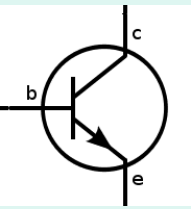


$$A(j\omega) = \frac{A_0 j \frac{\omega}{\omega_L}}{1 + j \frac{\omega}{\omega_L}}, \quad F = F_0 = \frac{R_1}{R_1 + R_2} < 1$$



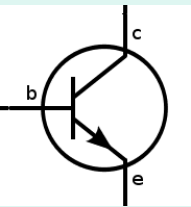
L08 Q03 HP amplifier with feedback

- What is the effect of the negative feedback on the critical frequency ω_L ?
 - A. It is not changed
 - B. It is increased
 - C. It is decreased



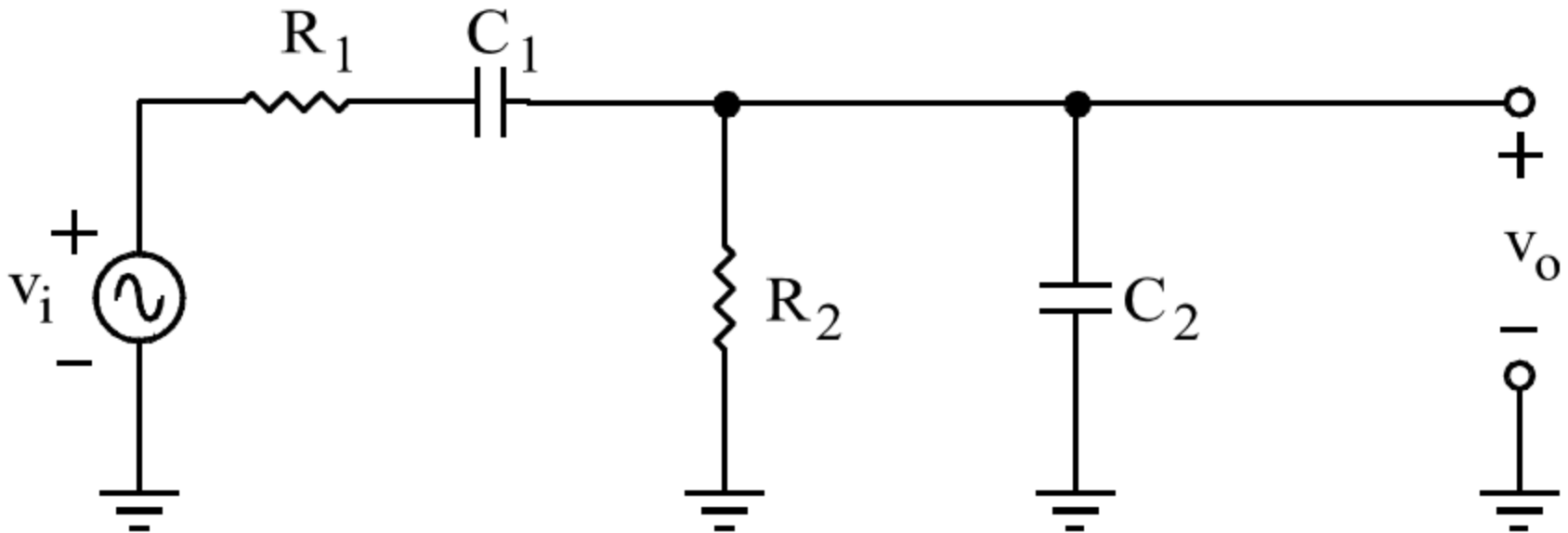
Open-circuit/short-circuit time constant method

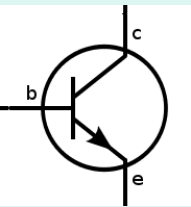
- Most circuits we operate with are bandpass, with capacitors as the only reactive components
- We desire a fast method to approximate the critical frequencies ω_L , ω_H



Example - simple bandpass filter

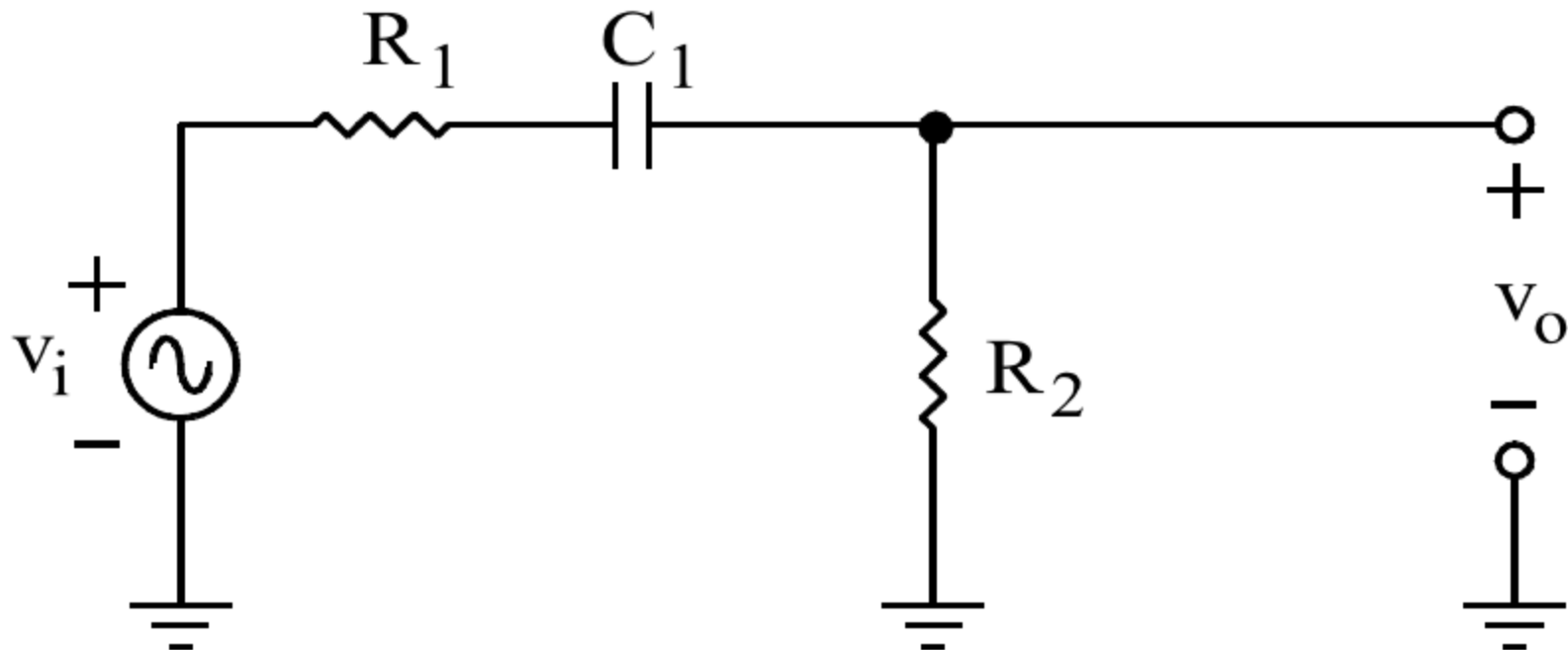
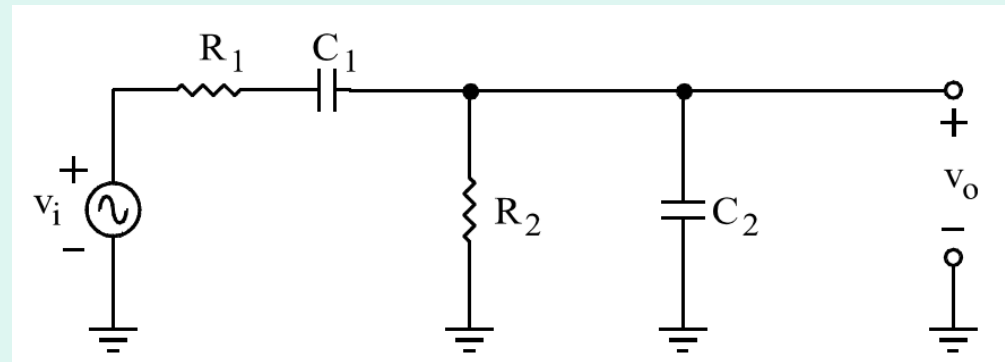
- C_1 blocks the LF components, C_2 attenuates the HF components
- To be a bandpass filter $C_1 \gg C_2$

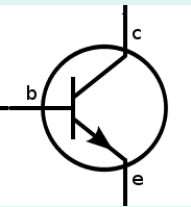




Circuit at LF

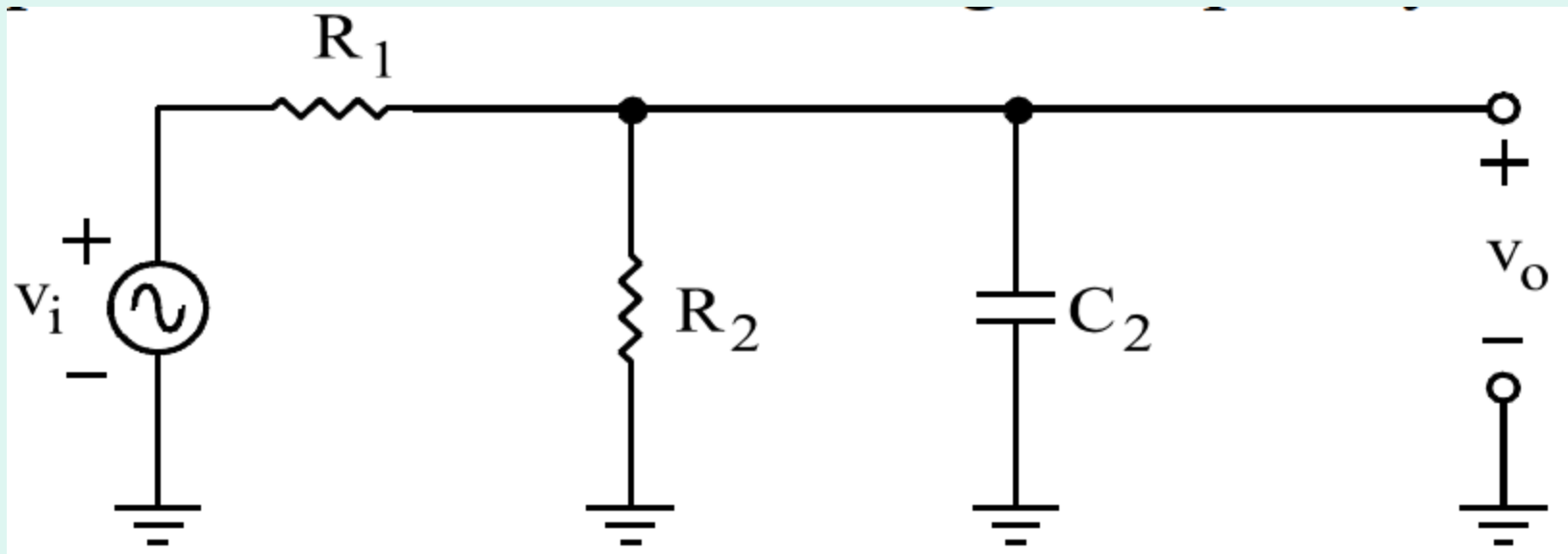
- Assume $C_1 \gg C_2$ - circuit behaves as high-pass filter for low frequency range

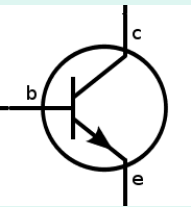




Circuit at HF

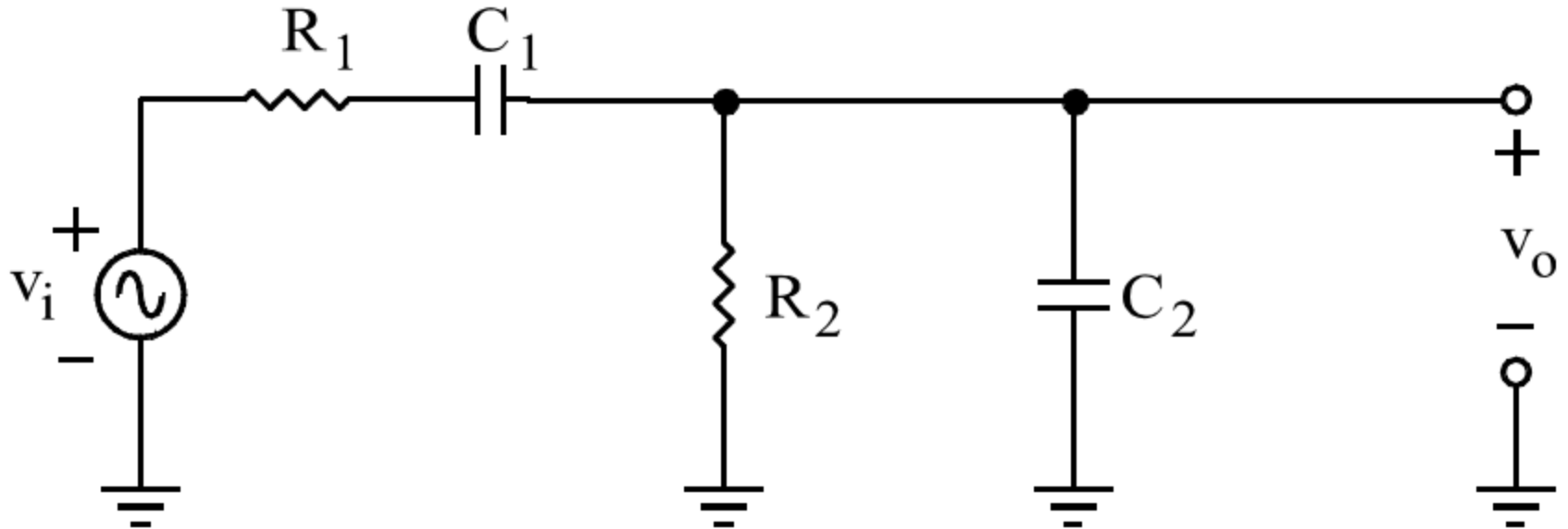
- Assume $C_1 \gg C_2$ - circuit behaves as low-pass filter for high frequency range



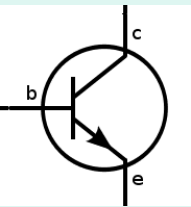


Transfer function

- Homework!



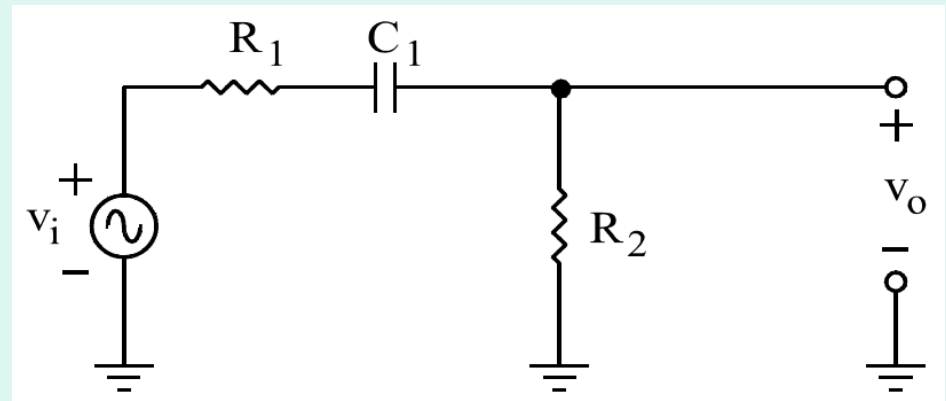
$$\frac{V_o(s)}{V_i(s)} = \frac{\frac{s}{R_1 C_2}}{s^2 + \left[\frac{1}{R_2 C_2} + \frac{1}{R_1 C_2} + \frac{1}{R_1 C_1} \right] s + \frac{1}{R_1 R_2 C_1 C_2}}$$



Approximate LF pole location

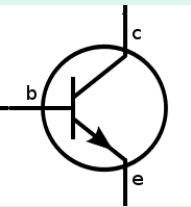
- Recall the circuit approximations for the LF and HF cases

LF regime



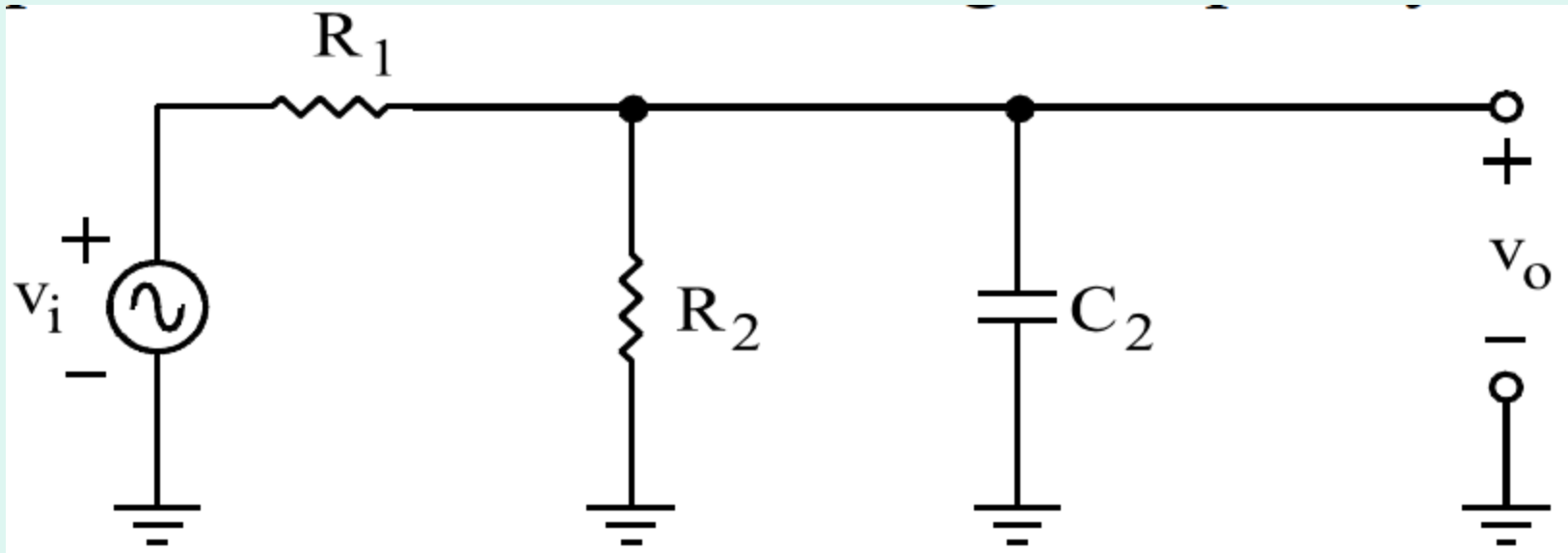
$$\left(\frac{V_o}{V_i} \right)_{LF} = \frac{R_2}{R_2 + R_1 + \frac{1}{sC_1}} = \frac{sR_2C_1}{1 + s(R_1 + R_2)C_1}$$

$$\omega_{p1} \approx \frac{1}{(R_1 + R_2)C_1} = \frac{1}{R_+ C_1}$$



Approximate HF pole location

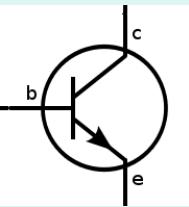
- The simplified circuit at HF:



$$\left(\frac{V_o}{V_i} \right)_{HF} \approx \frac{R_2 // \frac{1}{sC_2}}{R_1 + R_2 // \frac{1}{sC_2}} = \frac{\frac{R_2}{1 + sR_2C_2}}{R_1 + \frac{R_2}{1 + sR_2C_2}} = \frac{R_2}{R_1 + R_2 + sR_1R_2C_2}$$

$$\left(\frac{V_o}{V_i} \right)_{HF} \approx \frac{R_2}{R_1 + R_2} \frac{1}{1 + sC_2 \left(\frac{1}{R_1} + \frac{1}{R_2} \right)^{-1}}$$

$$\omega_{p2} \approx \frac{1}{\left(\frac{1}{R_1} + \frac{1}{R_2} \right)^{-1} C_2} = \frac{1}{(R_1 // R_2) C_2}$$



Combined frequency response

- With the obtained poles approximations:

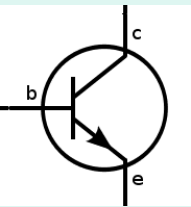
$$\frac{V_o(s)}{V_i(s)} = \frac{\frac{s}{R_1 C_2}}{s^2 + \left[\frac{1}{R_2 C_2} + \frac{1}{R_1 C_2} + \frac{1}{R_1 C_1} \right] s + \frac{1}{R_1 R_2 C_1 C_2}}$$

$$\frac{V_o(s)}{V_i(s)} = \frac{\frac{s}{R_1 C_2}}{\left(s + \frac{1}{R_+ C_1} \right) \left(s + \frac{1}{R_{||} C_2} \right)} = \frac{R_2}{R_1 + R_2} \frac{s}{\left(s + \frac{1}{R_+ C_1} \right)} \frac{\frac{1}{R_{||} C_2}}{\left(s + \frac{1}{R_{||} C_2} \right)}$$

$$= \frac{R_2}{R_1 + R_2} \frac{s}{(s + \omega_{p1})} \frac{\omega_{p2}}{(s + \omega_{p2})}$$

↗
↑
↖

Midband Lowfreq. Highfreq.



Generalization

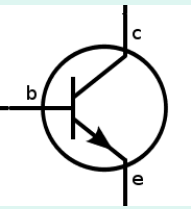
- We can in general, provided that the poles groups are apart, separate an amplifier frequency response into the three bands components:

$$T(s) = A_M F_L(s) F_H(s)$$

↗ ↑ ↖
Midband Lowf Highf
 (high-pass) (low-pass)

$$F_L(s) = \frac{(s + \omega_{Lz1}) \cdots (s + \omega_{Lzn})}{(s + \omega_{Lp1}) \cdots (s + \omega_{LpN})}$$

Here $n=N$ to ensure that $|F_L(\infty)|=1$



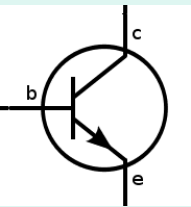
Generalization (2)

- For the HF component:

$$F_H(s) = \frac{\omega_{Hp1} \cdots \omega_{HpM}}{\omega_{Hz1} \cdots \omega_{Hzm}} \frac{(s + \omega_{Hz1}) \cdots (s + \omega_{Hzm})}{(s + \omega_{Hp1}) \cdots (s + \omega_{HpM})}$$

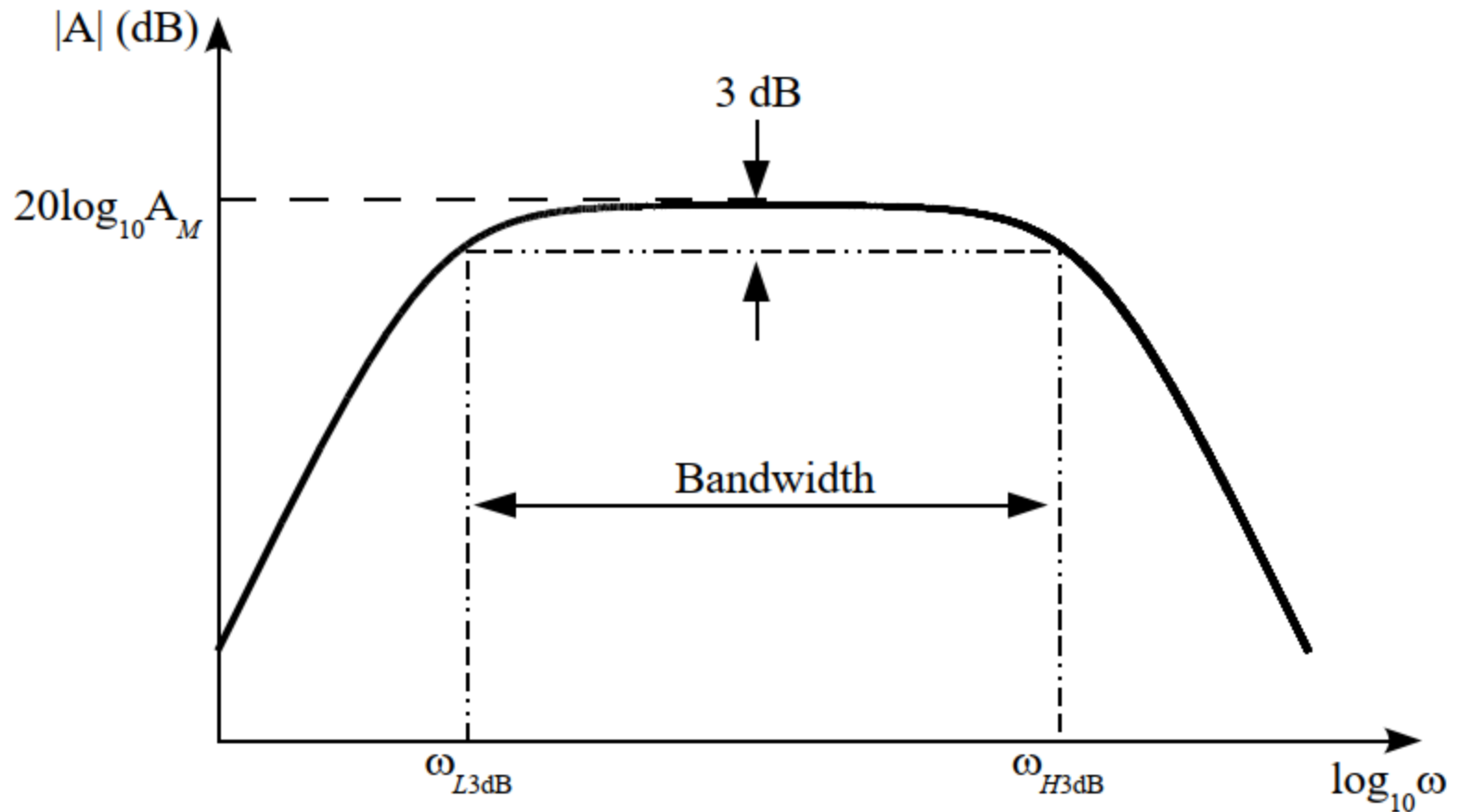
$$= \frac{\left(1 + \frac{s}{\omega_{Hz1}}\right) \cdots \left(1 + \frac{s}{\omega_{Hzm}}\right)}{\left(1 + \frac{s}{\omega_{Hp1}}\right) \cdots \left(1 + \frac{s}{\omega_{HpM}}\right)}$$

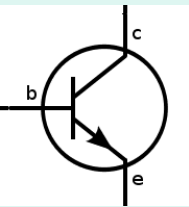
Here $M > m$, to ensure that $|F_H(\infty)| = 0$



Generic amplifier frequency response

- Bandwidth $BW = \omega_{H3dB} - \omega_{L3dB}$





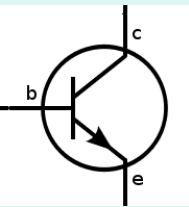
Finding the cut-off frequency ω_{L3dB}

- We assume we know the transfer functions (locations of poles and zeros)

$$|F_L(j\omega_{L3dB})|^2 = \frac{(\omega_{L3dB}^2 + \omega_{Lz1}^2)(\omega_{L3dB}^2 + \omega_{Lz2}^2) \cdots (\omega_{L3dB}^2 + \omega_{Lzn}^2)}{(\omega_{L3dB}^2 + \omega_{Lp1}^2)(\omega_{L3dB}^2 + \omega_{Lp2}^2) \cdots (\omega_{L3dB}^2 + \omega_{LpN}^2)} = \frac{1}{2}$$

$$2(\omega_{L3dB}^2 + \omega_{Lz1}^2) \cdots (\omega_{L3dB}^2 + \omega_{Lzn}^2) = (\omega_{L3dB}^2 + \omega_{Lp1}^2) \cdots (\omega_{L3dB}^2 + \omega_{LpN}^2)$$

$$\begin{aligned} & 2(\omega_{L3dB}^{2n} + (\omega_{Lz1}^2 + \omega_{Lz2}^2 + \cdots + \omega_{Lzn}^2)\omega_{L3dB}^{2(n-1)} + \cdots + (\omega_{Lz1}^2 \omega_{Lz2}^2 \cdots \omega_{Lzn}^2)) \\ &= (\omega_{L3dB}^{2N} + (\omega_{Lp1}^2 + \omega_{Lp2}^2 + \cdots + \omega_{LpN}^2)\omega_{L3dB}^{2(N-1)} + \cdots + (\omega_{Lp1}^2 \omega_{Lp2}^2 \cdots \omega_{LpN}^2)) \end{aligned}$$

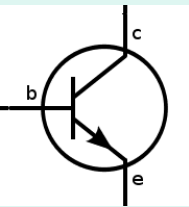


Finding the cut-off frequency ω_{L3dB} (2)

- Since ω_{L3dB} is larger than any pole or zero, we can approximate the eqn. considering only the highest powers
- Rem: many of the zeros may be located at 0

$$2\omega_{L3dB}^{2N} + 2(\omega_{Lz1}^2 + \dots + \omega_{Lzn}^2)\omega_{L3dB}^{2(N-1)} \approx \omega_{L3dB}^{2N} + (\omega_{Lp1}^2 + \dots + \omega_{LpN}^2)\omega_{L3dB}^{2(N-1)}$$

$$\omega_{L3dB} \approx \sqrt{\omega_{Lp1}^2 + \dots + \omega_{LpN}^2 - 2\omega_{Lz1}^2 - \dots - 2\omega_{Lzn}^2}$$

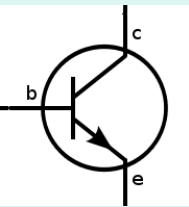


Finding the cut-off frequency ω_{3dBH}

- At ω_{3dBH} , the normalized $F_H(s)$ should be reduced by $\sqrt{2}$

$$|F_H(j\omega_{H3dB})|^2 = \frac{\left(1 + \frac{\omega_{H3dB}^2}{\omega_{Hz1}^2}\right) \left(1 + \frac{\omega_{H3dB}^2}{\omega_{Hz2}^2}\right) \dots \left(1 + \frac{\omega_{H3dB}^2}{\omega_{Hzm}^2}\right)}{\left(1 + \frac{\omega_{H3dB}^2}{\omega_{Hp1}^2}\right) \left(1 + \frac{\omega_{H3dB}^2}{\omega_{Hp2}^2}\right) \dots \left(1 + \frac{\omega_{H3dB}^2}{\omega_{HpM}^2}\right)} = \frac{1}{2}$$

$$2 \left(1 + \frac{\omega_{H3dB}^2}{\omega_{Hz1}^2}\right) \left(1 + \frac{\omega_{H3dB}^2}{\omega_{Hz2}^2}\right) \dots \left(1 + \frac{\omega_{H3dB}^2}{\omega_{Hzm}^2}\right) = \left(1 + \frac{\omega_{H3dB}^2}{\omega_{Hp1}^2}\right) \left(1 + \frac{\omega_{H3dB}^2}{\omega_{Hp2}^2}\right) \dots \left(1 + \frac{\omega_{H3dB}^2}{\omega_{HpM}^2}\right)$$



Finding the cut-off frequency $\omega_{3\text{dBH}}$ (2)

- Since $M > m$ and $\omega_{3\text{dBH}}$ is smaller than any of the pole or zero frequencies, we can approximate the equality considering only the (dominant) low powers in the polynomials

$$2 + \left(\frac{2}{\omega_{Hz1}^2} + \dots + \frac{2}{\omega_{Hzm}^2} \right) \omega_{H3\text{dB}}^2 \approx 1 + \left(\frac{1}{\omega_{Hp1}^2} + \dots + \frac{1}{\omega_{HpM}^2} \right) \omega_{H3\text{dB}}^2$$

$$\frac{1}{\omega_{H3\text{dB}}^2} \approx \frac{1}{\omega_{Hp1}^2} + \dots + \frac{1}{\omega_{HpM}^2} - \frac{2}{\omega_{Hz1}^2} - \dots - \frac{2}{\omega_{Hzm}^2}$$

$$\tau_{H3\text{dB}} \approx \sqrt{\tau_{Hp1}^2 + \dots + \tau_{HpM}^2 - 2\tau_{Hz1}^2 - \dots - 2\tau_{Hzm}^2}$$

$$\omega_{H3\text{dB}} = \frac{1}{\tau_{H3\text{dB}}}$$

